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AERODYNAMIC CHARACTERISTICS OF

THE 10-PERCENT-THICK NASA

SUPERCRITICAL AIRFOIL 33 DESIGNED FOR A

NORMAL-FORCE COEFFICIENT OF 0.7

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AERODYNAMIC CHARACTERISTICS OF

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By Charles D. Harris

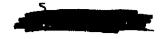
Langley Research Center

SUMMARY

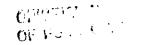
A 10-percent-thick supercritical airfoil based on an off-design sonicpressure plateau criterion has been developed and experimental aerodynamic
characteristics measured. The airfoil had a design normal-force coefficient
of 0.7 and was identifie⁴ as supercritical airfoil 33. Results show the airfoil to have good drag rise characteristics over a wide range of normal-force
coefficients with no measurable shock losses up to the Mach numbers at which
drag divergence occurred for normal-force coefficients up to 0.7. Comparisons
of experimental and theoretical characteristics were made and composite drag
rise characteristics were derived for normal-force coefficients of 0.5 and 0.7
and a Reynolds number of 40 million.

INTRODUCTION

Continued development of supercritical airfoil technology has resulted in recognition of design criteria which permit the design of family related supercritical airfoils. Based on these criteria, two supercritical airfoils have been designed - the 10-percent thick airfoil 33 reported herein and the 14-percent-thick airfoil of reference 1. The design normal-force coefficient was 0.7 for both airfoils.



SYMBOLS



Values are given in both SI and U.S. Customary Units. Measurements and calculations are made in U.S. Customary Units.

C _p	pressure	coefficient,	<u>p</u> - p	<u>.</u>
			\mathtt{d}^{∞}	

Cp, sonic pressure coefficient corresponding to local Mach number of 1.0

c chord of airfoil, 63.5 centimeters (25.0 inches)

c_d section drag coefficient

Acd,s drag increment due to shock wave losses

e section lift coefficient

section pitching-moment coefficient about the quarter-chord point

 $c_{f n}$ section normal-force coefficient

K surface curvature, reciprocal of local radius of curvature

M Mach number

m surface slope, dy/dx

p static pressure, newtons per meter² (pounds per foot²)

Δp_t total-pressure loss, newtons per meter² (pounds per foot²)

q dynamic pressure, newtons per meter² (pounds per foot²)

R Reynolds number based on airfoil chord

x ordinate along airfoil reference line measured from airfoil leading edge, centimeters (inches)

y ordinate normal to airfoil reference line, centimeters (inches)

vertical distance in wake profile measured from bottom of rake, centimeters (inches)

a geometric angle of attack of airfoil reference line, degrees

2



Subscripts:

- l local point on airfoil
- " undisturbed stream

APPARATUS AND TECHNIQUES

Model Configuration

The supercritical airfoil basic concept and detailed design philosophy are discussed in reference 2.

Background.- During initial phases of the two-dimensional supercritical airfoil development program, emphasis was placed upon developing a 10-percent-thick airfoil with the highest drag divergence Mach number attainable at a normal-force coefficient of 0.7. The normal-force coefficient of 0.7 was chosen as the design goal since, when account was taken of the sweep effect, it was representative of lift coefficients at which the advanced technology near-sonic transports utilizing the supercritical airfoil concept were then expected to cruise. These initial phases resulted in supercritical airfoil 11 which has been reported in reference 3. This airfoil exhibited an undesirable creep or gradual increase in the variation of drag coefficient with Mach number of about 14 counts (c_d increment of 0.0014) between M = 0.50 and the drag divergent Mach number at the design normal-force coefficient.

Subsequent design studies of advanced technology transport configurations suggested that the near-sonic cruise lift coefficient requirements would be somewhat lower than originally anticipated. Consequently, the improved supercritical airfoil 26a (ref. 4) was developed with a design normal-force coefficient of about 0.55. The wind-tunnel tests required for this airfoil also provided the opportunity to explore the drag creep problem noted with airfoil 11. Airfoil 26a had no drag creep at normal-force coefficients up to about 0.6 and



the drag divergence Mach number varied from approximately 0.82 at a normalforce coefficient of 0.30 to 0.78 at a normal-force coefficient of 0.80. As discussed in reference 4 these improved drag creep characteristics were largely, attributed to a more favorable flow recompression over the forward upper surface and the elimination of a region of overexpansion near the three-quarterchord station.

Recent emphasis on fuel economy has generated considerable interest in a fuel-conserving aircraft envisioned to cruise at Mach numbers near those of current transports. Such an aircraft could utilize supercritical airfoil technology to achieve weight and drag reductions by permitting the use of thicker wings with higher aspect ratios and less sweep. Because the wings with higher aspect ratio would require airfoils with design normal-force coefficients higher than 0.55, airfoil 31 (ref. 5) was developed with the same design normal-force coefficient as the earlier airfoil 11 ($c_{
m n}$ = 0.7) but with significant improvements in the drag characteristics. Drag creep was practically eliminated at normal-force coefficients between about 0.4 and 0.7 and greatly reduced at other normal-force coefficients. Substantial reductions in the drag levels preceding drag divergence were also achieved at all normal-force coeffi-

Airfoils through 31 were developed experimentally through intuitive contour modifications in the wind tunnel. It was an iterative process consisting of evaluating experimental pressure distributions at design and off-design conditions and physically altering the airfoil profile to yield the best drag characteristics over a range of test conditions.

The experimental results for these earlier airfoils suggested design criteria which would yield the best compromise in drag characteristics over a



range of test conditions. These design criteria consist essentially of three principal guidelines which may be used in designing supercritical airfoils. There are several additional, more detailed, design guidelines (treatment of leading and trailing edges, and local minimum thickness constraints, for example) which are beyond the intended scope of this report and will be discussed in a later report.

The first principal guideline, referred to as the sonic plateau criterion, is that at some incremental normal-force coefficient below the design normal-force coefficient the pressure distribution on the upper and lower surfaces be flat with the upper surface pressures just below the sonic value. The increment is a function of the design normal-force coefficient and appears to be about -0.25 to -0.30 for a design normal-force coefficient of 0.7.

On the upper surface the plateau extends from near the leading edge to the start of the aft pressure recovery and on the lower surface from near the leading edge to the recompression region entering into the cusp. The rearward extent of the upper surface plateau is determined by the second principal design guideline which requires that the gradient of the aft pressure recovery be gradual enough to avoid separation problems for lift coefficients and Mach numbers up to the design value. Consequently, the rearward extent of the upper surface plateau would depend on thickness ratio since the thicker the airfoil, the higher the induced velocities from which the flow must recover and, therefore, the further forward the aft pressure recovery must begin. The upper surface plateau extends from approximately 3- to 80-percent chord on the 10-percent-thick airfoil 33 and approximately 5- to 66-percent chord on the 14-percent-thick airfoil of reference 1.

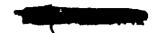


The third principal guideline requires that the airfoil have sufficient aft camber so that at design conditions the angle of attack be about zero. This prevents a too-forward location of the upper surface crest with the negative pressure coefficients over the mid-chord region acting over a rearward facing surface. Both experiments and theoretical analyses (ref. 6) have indicated that an increase in angle of attack to positive values results in an abrupt increase in wave drag.

Based on these criteria two supercritical airfoils were designed - the 14-percent-thick airfoil reported in reference 1 and the 10-percent-thick airfoil 33 reported herein. The design normal-force coefficient was 0.7 for both airfoils. An iterative design process was used which consisted of altering the airfoil coordinates until the viscous, airfoil analysis program of reference 6 indicated that the aforementioned design criteria had been satisfied. Geometric characteristics of the experimental 10-percent-thick airfoil 33 are presented in figures 1 and 2 and compared with those of the earlier 10-percent-thick airfoil 31 (ref. 5). Measured section coordinates of both airfoils are presented in table I.

The incremental ordinates between airfoils 31 and 33 delineated in table I show that the ordinates were modified over the forward upper and lower surfaces, decreased over the rear upper surface, and increased in the vicinity of the 80-percent chord on the lower surface. Referring to the pressure distribution for airfoil 31 which approaches the off-design sonic-plateau criterion (fig. 15(c) of ref. 5; M = 0.76 and $c_n = 0.46$), the alterations over the upper surface and forward lower surface were necessary to obtain the desired plateau pressure distributions and to reduce the upper surface aft pressure recovery gradient.

The ordinates on the rear lower surface were increased, with the maximum increase



at 80-percent, to provide increased depth for control surface and flap structural requirements. Subtracting from the upper surface and adding to the lower surface over the aft portion of the airfoil in this manner reduced the aft camber and, therefore, increased the angle of attack required to achieve a given normal-force coefficient. The ordinates on the lower surface were progressively increased until the third design criterion was met; that is, that the angle of attack required for the design normal-force coefficient of 0.7 was approximately 0°.

Since the best drag characteristics are often obtained on airfoils with a small amount of upper surface trailing-edge separation and since theoretical treatments of the flow at trailing-edge regions are generally unreliable, theoretically-predicted flow separation at the 98-percent chord location was accepted during the design process. Attempts to achieve a more rearward location of theoretical separation by reducing the aft pressure recovery gradient would have forced the rear terminus of the sonic plateau forward, resulting in higher induced velocities in the plateau region and a probable reduction in drag rise Mach number.

Wind-tunnel model. The coordinates of the experimental airfoil 33 deviated slightly from the design profile (not presented). On the upper surface, the deviation was nowhere greater than $\Delta x/c = 0.0003$ and generally less than 0.0001. On the lower surface, deviation was somewhat greater; as much as 0.0005 on the forward region. These small deviations should not significantly affect the results.

Irregularities in the curvature distributions (fig. 2) are due to small surface irregularities that become greatly exaggerated when examined from the standpoint of local curvature. Such irregularities are not as apparent in the slope distributions (fig. 2). Both airfolds included a trailing-edge cavity



(see the insert in fig. 1 and the photographs of fig. 3) which had a favorable effect on the wake as discussed in reference 3.

The wind-tunnel models, mounted in an inverted position, spanned the width of the tunnel with a span-chord ratio of 3.43. They were constructed with metal leading and trailing edges and a metal core around which plastic fill was used to form the contours of the airfoils. Angle of attack was changed manually by rotating the model about pivots in the tunnel sidewalls. A photograph and a drawing of a typical airfoil model installed in the tunnel are shown in figures 3 and 4, respectively.

Wind Tunnel

The investigation was conducted in the Langley 8-foot transonic pressure tunnel (ref. 7). This tunnel is a continuous-flow, variable-pressure wind tunnel with controls that permit the independent variation of Mach number, stagnation pressure and temperature, and dewpoint. It has a 2.16-meter-square (85.2-inch-square) test section with filleted corners so that the total cross-sectional area is equivalent to that of a 2.44-meter-diameter (8-foot-diameter) circle. The upper and lower test-section walls are axially slotted to permit testing through the transonic speed range. The total slot width at the position of the model averaged about 5 percent of the width of the upper and lower walls.

The solid sidewalls and slotted upper and lower walls make this tunnel well suited to the investigation of two-dimensional models since the sidewalls act as endplates and the slots permit development of the flow field in the vertical direction.

Boundary-Layer Transition

Based on the technique discussed in reference 8, boundary-layer transitio. was fixed along the 28-percent chord line on the upper and lower surfaces of

the models in an attempt to simulate full-scale Reynolds numbers by providing the same relative trailing-edge boundary-layer displacement thickness at model scale as would exist at full-scale flight conditions. The simulation technique, which requires that laminar flow be maintained ahead of the transition trip, is limited to those test conditions in which shock waves or steep adverse pressure gradients occur behind the point of fixed transition so that the flow is not tripped prematurely. The transition trips consisted of 0.25-cm-wide (0.10-in.) bands of No. 90 Carborundum grains. Analysis of drag and boundary-layer displacement thickness at the trailing edge theoretically computed using reference 6 indicated that the simulated full-scale Reynolds number was around 40 million rather than the 20 to 30 million quoted for earlier supercritical airfoil investigations (ref. 3, for example).

Measurements

Surface-pressure measurements. Normal force and pitching moments on the airfoils were determined from surface static-pressure measurements. The surface-pressure measurements were obtained from a chordwise row of orifices located approximately 0.32c from the tunnel center line. Orifices were more concentrated near the leading and trailing edges of the airfoil to define the pressure gradients in these regions. In addition, a rearward facing orifice was included in the cavity at the trailing edge (identified at an upper surface x/c location of 1.00). The transducers used in the differential pressure scanning valves to measure the static pressure at the airfoil surface had a range of ±68.9 kN/m² (10 1b/in²).

Wake measurements. - Drag forces were determined from vertical variations
of the total and static pressures measured across the wake with the profile drag
rake shown in figure 4(b). The profiles, schematically illustrated in figure 5,
represent the momentum losses as indicated by stagnation-pressure deficits

across the wake. The middle section of these profiles reflects viscous and separation losses in the boundary layer, whereas the "wings" of the profile reflect direct losses in stagnation pressure across the shock waves.

The rake was positioned in the vertical center-line plane of the tunnel, approximately 1 chord length rearward of the trailing edge of the airfoil. The total-pressure tubes were flattened horizontally and closely spaced vertically (0.36 percent of the airfoil chord) in the region of the wake associated with skin-friction boundary-layer losses. Outside this region, the tube vertical spacing progressively widened until in the region above the wing where only shock losses were anticipated, the total-pressure tubes were spaced apart about 7.2 percent of the chord. Static-pressure tubes were distributed as shown in figure 4(b). Each static pressure measured was used over a section of the rake to determine local flow conditions in the vicinity of the static-pressure tube rather than using an average of all the static pressures measured. The rake was attached to the conventional center line sting mount of the tunnel; this arrangement permitted it to be moved vertically to center the close concentration of tubes in the boundary-layer wake. The transducer in the differentialpressure scanning valve connected to total-pressure tubes intended to measure boundary-layer losses had a range of $\pm 17.2 \text{ kN/m}^2$ (2.5 lb/in²), and the transducers in the valves for measuring shock losses and static pressure had a range of $\pm 6.9 \text{ kN/m}^2 (1 \text{ lb/in}^2)$.

Reduction of Data

Calculation of c_n and c_m . Section normal-force and pitching-moment coefficients were obtained by numerical integration (based on the trapezoidal method) of the local surface-pressure coefficient measured at each orifice multiplied by an appropriate weighting factor (incremental area).

Calculation of c_d . To obtain section drug coefficients, point drag coefficients were computed for each total-pressure measurement in the wake by using the procedure of reference 9. These point drag coefficients were then summed by numerical integration across the wake, again based on the trapezoidal method. Drag increments due to shock wave losses ($\Delta c_{d,s}$) were determined from integration of the drag measured across the wings (fig. 5) of the wake profile.

Wind-Tunnel-Wall Effects

Because of the uncertainty in lift-induced interference effects and solid and wake blockage effects (particularly in the presence of local supercritical flow) no corrections for wall effects have been applied to the basic experimental data. Adjustments for blockage were applied to the composite drag rise data and are explained in the Discussion section.

TEST CONDITIONS

Tests were conducted at Mach numbers from 0.50 to 0.82 for a stagnation pressure of 0.1013 MN/m² (1 atm). The stagnation temperature of the tunnel air was automatically controlled at approximately 322 K (120°F) and the air was dried until the dewpoint in the test section was reduced sufficiently to avoid condensation effects. Resultant test Reynolds numbers based on the airfoil chord length were as shown in figure 6.

PRESENTATION OF RESULTS

The experimental data reported herein are presented in the following figures:

		F	igure
ų	Force and moment characteristics		7
	Variation of measured section drag coefficient with Mach number	•	8
•	Drag increment due to shock-wave losses		Q

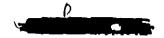


	Figure
Composite drag curves	. 10
Chordwise pressure distributions at -	,
M = 0.50	. 11
$M = 0.60. \dots$. 12
M = 0.70	. 13
M = 0.74	. 14
M = 0.76	. 15
$M = 0.77. \dots$. 16
M = 0.78.	. 17
M = 0.79	. 18
M = 0.80	. 19
M = 0.81	. 20
M = 0.82.	- 21
Comparisons of theoretical and experimental characteristics for -	
$M = 0.76$, $c_n = 0.414$. 22
$M = 0.79, c_n = 0.535$	23
$M = 0.77$, $c_n = 0.707$. 24
Complete surface pressure distributions for the earlier 10-per	ent-thick
supercritical airfoil 31 are presented in reference 5 and for the fa	
3h	

14-percent-thick airfoil in reference 1.

DISCUSSION

Measured Aerodynamic Characteristics

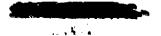
Sonic plateau. - Figure 15(c) indicates how close the analytically designed, experimental 10-percent-thick airfoil 33 came to satisfying the off-design sonic, plateau criterion for M = 0.76 and $c_n = 0.41$. Irregularities in the plateau



pressures are due to small surface imperfections which become greatly exaggerated when the flow is right on the verge of sonic velocity. The upper surface plateau extends from approximately 3- to 80-percent chord and comes much nearer to satisfying the sonic-plateau criterion than that for airfoil 31 (ref. 5). The aft pressure recovery gradient is considerably reduced and the trailing-edge pressure recovery more positive, compared with airfoil 31.

The characteristic supercritical pressure distribution at near design conditions (slightly decelerating upper surface velocities terminated by a very weak recompression near the midchord and followed by a near-sonic plateau before entering the final trailing-edge pressure recovery) for an angle of attack near the design value of 0° (fig. 24) occurs at the experimental conditions of M = 0.77 and $c_n = 0.71$ (fig. 16(f)). The increment between normal-force coefficients at design and off-design sonic plateau conditions is, therefore, about -0.25 to -0.30 as suggested.

Measured drag characteristics. Figure 8 indicates that although airfoil 33 had slightly lower drag divergence Mach numbers than airfoil 31 because of earlier onset of shock losses (fig. 9), the drag characteristics of wirfoil 33 preceding drag divergence were much improved compared to those of airfoil 31. The improved drag characteristics preceding drag divergence are associated primarily with reduced aft adverse pressure gradients which result in less profile drag. The increased wave losses are associated with the more positive angle of attack required to achieve the same normal-force coefficients as airfoil 31 because of the reduced aft camber and with the influence of the increased curvature over the upper surface mid-chord region (fig. 2(a)). These increased wave losses occur at conditions beyond the design conditions however, and were stolerated in exchange for enhanced performance preceding drag divergence.



The apparent dips in the drag rise curves (fig. 8) should not be interpreted as being due to reductions in shock losses as the design Mach number is approached. At $c_n = 0.7$, for example, there is a dip near M = 0.77 but there are no shock losses evident (fig. 9) below M = 0.77. In an attempt to simulate full scale Reynolds numbers for the design Mach number, boundary-layer transition trips were fixed at 28-percent chord. At M = 0.70, the pressure distribution for $c_n = 0.7$ (fig. 13) was such that laminar flow could not be maintained back to the trip. The higher drag level preceding M = 0.77 was, therefore, due to premature boundary-layer transition with greater skin friction drag than there would have been if laminar flow could have been maintained back to the trip at 28-percent chord. Higher drag levels preceding the dip are not associated with increased form drag due to separation effects since no losses in trailing-edge pressure recovery were evident.

Measured pitching-moment characteristics.— The negative or leading edge down pitching moments of airfoil 31 (fig. 7) were reduced by about $\Delta c_m = 0.01$ because of the reduced aft camber incorporated into airfoil 33. At design conditions (M \approx 0.78 and $c_n = 0.7$ for airfoil 33 (fig. 7), and M \approx 0.74 and $c_n = 0.7$ for the 14-percent-thick airfoil of reference 1) pitching-moment coefficients were practically the same for the family related 10- and 14-percent airfoils.

Comparison With Theory

Correlation was established between experimental and theoretically predicted data using the viscous, analysis program of reference 6. Representative results are presented in figures 22 to 24. Since the viscous analysis program was known to overpredict trailing-edge pressure recovery, Mach number was varied, to achieve the best matches of the theoretical and experimental pressure



distributions over the forward regions of the airfoil and shock wave locations for given normal-force coefficients. The results indicate the Mach number and angle of attack adjustments needed to correct for tunnel wall interference effects.

When the flow over the model is subcritical (fig. 22, for example) the blockage or Mach number correction required to match the experimental and theoretical pressure distributions over the forward region of the airfoil are small and tend to agree with what would be predicted by subcritical theory (ref. 10, for example). When substantial amounts of supercritical flow begin to appear over the airfoil blockage corrections become significant as indicated by figures 23 and 24.

Attention must be called to two important points in the correlations shown in figures 22 to 24. The total theoretical drag calculated by the viscous, analysis program of reference 6 is calculated in two parts; profile drag, consisting of skin friction and form or pressure drag, and the contribution of wave drag which exists in supersonic flow. Experience has indicated that because of overprediction of wave losses the total theoretical drag tends to become too large as soon as supercritical flow appears on the airfoil. Consequently, the theoretical drag shown in figures 22 to 24 is that drag associated with only profile drag and agrees well with the experimental drag. If the wave losses as calculated by the method of reference 6 were taken into account, the theoretical drag would be greater than the experimental drag by 0.0010 and 0.0019 for figures 23 and 24 respectively.

The second point which must be made concerning the correlation figures pertains to the position of boundary-layer transition specified for the theoretical data. Reference 6 neglects the laminar portion of the flow ahead of





the transition point. Because of the lengthy run of laminar flow (28-percent) and its effect on the turbulent boundary-layer development on the model, the neglected laminar boundary had to be taken into account. This was accomplished by specifying a transition point forward of 28-percent chord in the theoretical calculations to get better agreement between the theoretical and experimental drag for flow conditions with no supercritical flow (zero wave losses) and where the experimental transition occurred at 28-percent chord. The theoretical calculations shown in figures 22 to 24 were, therefore, calculated for tunnel Reynolds numbers with transition fixed at 24-percent chord.

Composite Drag and Angle-of-Attack Characteristics

Drag. - A combination of experimental and theoretical drag characteristics for airfoil 33 is shown in figure 10 in order to synthesize a realistic drag rise curve for a full-scale Reynolds number. Theoretical drag values (solid line) based on 40 million Reynolds number and boundary-layer transition at 3-percent chord are used for Mach numbers up to the Mach number at which shock losses become evident in the experimental data. Because of the inaccuracy of the theoretical wave losses, experimental drag values (dashed line) are used beyond that point. The experimental and theoretical drag agree at the junction of the solid and dashed lines. The nominal Mach numbers for the experimental data were reduced by the increments indicated by figures 23 and 24 to be required to account for blockage effects.

The gradual increase in drag with Mach number up to drag divergence was associated with increased profile drag due to the effect of Mach number on the induced velocities. As noted in an earlier section there were no measurable shock losses up to the Mach numbers at which drag diverged.

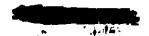
Angle of attack.- Theoretical angles of attack required to obtain the desired section lift coefficients are also shown in figure 10 and are near zero at the design or cruise Mach numbers. The difference in angles of attack for $c_n=0.5$ and 0.7 provide an indication of the lift curve slope in this lift range.

CONCLUDING REMARKS

A 10-percent-thick supercritical airfoil based on an off-design sonic pressure plateau criterion has been developed and experimental aerodynamic characteristics measured. The airfoil had good drag rise characteristics over a wide range of normal-force coefficients with no measurable shock losses up to the Mach numbers at which drag divergence occurred for normal-force coefficients up to 0.7. Comparisons of experimental and theoretical characteristics were made and composite drag rise characteristics were derived for normal-force coefficients of 0.5 and 0.7 and a Reynolds number of 40 million.



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TABLE I.- SECTION COORDINATES

[c = 63.5 cm (25 in.); airfoil 31 has an uppersurface leading-edge radius of 0.0169c and a lower surface leading-edge radius of 0.0176c; airfoil 33 has a leading-edge radius of 0.015c; Δ(y/c) = (y/c) airfoil 33 - (y/c) airfoil 31]

x/c	(y/c) _u	(y/c) ₁	(y/c) _u	(y/c) ₁		
	Airfo	oil 31	Airfoil 33		Δ(y/c) _u	Δ(y/c) ₁
0.0	0.0	0.0	0.0	0.0	0.0	0.0
.002	.0080	0081	.0075	0075	0005	.0006
.005	.0120	0123	.0116	0116	0004	.0007
.010	.0159	0164	.0156	0156	0003	.0008
.020	.0209	0215	.0206	0203	0003	.0012
.030	.0244	0250	.0240	0235	0004	.0015
.040	.0271	0274	.0267	0262	0004	.0012
.050	.0294	0295	.0289	0284	0005	.0012
.060	.0313	0314	.0308	0303	0005	.0011
.070	.0329	0330	.0325	0320	0004	.0010
.080	.0344	0344	.0340	0335	0004	Ì
.090	.0357	0358	.0354	0349	0003	.0009
.100	.0369	0370	.0367	0362	0003	.0009
.110	.0380	0381	.0379	0374	0002	
.120	.0391	0392	.0389	0385	0002	.0007
.130	.0401	0402	.0399	0395	0002	.0007
.140	.0410	0411	.0408	0405	0002	.0007
.150	.0418	0420	.0417	0414	į.	.0006
.160	.0426	0427	.0425	0422	0001	.0006
.170	.0433	0434	.0432	0422	0001	.0005
.180	.0440	0440	.0432	0430	0001	.0004
.190	.0446	0446	.0439	0437	0001. 0. 0	.0003 .0002

TABLE I.- SECTION COORDINATES - Continued

x/c	(y/c) _u	(y/c) ₁	(y/c) _u	(y/c) ₁	11(.)	.,,,
,	Airfoi	1 31	Airfoi	1 33	Δ(y/c) _u	Δ(y/c) ₁
0.200	0.0452	-0.0452	0.0452	-0.0450	0.0	0.0002
.210	.0457	0457	.0457	0456	0.0	.0001
.220	.0462	0462	.0462	0462	0.0	0.0
.230	.0467	0467	.0467	0467	0.0	0 0
.240	.0471	0471	.0472	0472	.0001	0001
.250	.0475	0475	.0476	0476	.0001	0001
.260	.0478	0479	.0480	0480	.0002	0001
.270	.0481	0482	.0483	0484	.0002	0002
.280	.0484	0485	.0486	0487	.0002	0002
.290	.0487	0488	.0489	0490	.0002	0002
.300	.0489	0491	.0491	0493	.0002	0002
.310	.0491	0493	.0493	0495	.0002	0002
.320	.0493	0495	.0495	0497	.0002	0002
•330	.0495	0497	.0496	0499	.0001	0002
•340	.0497	0498	.0497	0500	0.0	0002
.350	.0498	0499	.0498	0501	0.0	0002
.360	.0499	0499	.0499	0502	0.0	0003
.370	.0500	0499	.0500	0502	0.0	0003
.380	.0500	0499	.0500	0502	0.0	0003
.390	.0500	0498	.0500	0502	0.0	0004
-400	.0500	0497	.0500	0501	0.0	0004
.410	.0500	0495	.0499	0500	0001	0005
.420	.0500	0493	.0498	0499	0002	0006
.430	.0499	0491	.0497	0497	0002	0006
.440	.0498	0488	.0496	0495	0002	0007
.450	.0497	0485	.0495	0492	0002	0007
.460	.0495	0482	.0493	0488	0002	0006
•470	.0493	0478	.0491	0484	0002	0006

TABLE I.- SECTION COORDINATES - Continued

x/c	(y/c) _u	(y/c) ₁	(y/e) _u	(y/c) ₁	Δ(y/c) _u	
	Airfo	oil 31	Airfo	Airfoil 33		Δ(y/c) ₁
0.48¢	0.0491	-0.0474	0.0489	-0.0480	-0.0002	-0,0006
•490	.0489	0470	-0487	0475	0002	0005
•500	.0487	0465	.0484	0470	0003	0005
•510	.0485	0459	.0481	0464	0004	0005
•520	.0482	0453	.0478	0457	0004	0004
•530	.0479	0446	.0475	0450	0004	0004
•540	.0476	0439	.0472	0442	0004	0003
•550	.0473	0431	.0468	0434	0005	0003
.560	.0469	0422	.0464	0425	0()5	0003
-570	.0465	0413	.0460	0415	0005	0003
.580	,0461	0403	.0456	0405	0005	0002
•590	.0457	0392	.0451	0394	0006	0002
.600	.0453	0381	.0446	0382	0007	0001
.610	.0448	0369	.0441	0370	0007	0001
.620	.0443	0356	.0436	0357	0007	0001
•630	.0438	0342	.0430	0343	0008	0001
.640	.0433	0327	.0424	0329	0009	0002
.650	.0428	0311	.0418	0315	0010	0004
.660	.0422	0294	.0412	0300	00T0	0006
.670	.0416	0277	.0405	0285	0011	0003
.680	.0409	0260	.0398	0270	0011	0010
.690	.0402	0242	.0391	0255	0011	013
.700	.0395	0224	•0383	0239	0012	0015
.710	.0387	0206	•0375	0223	0012	0017
•720	.0379	0188	.0367	0207	0012	0019
•730	.0371	0171	.0358	0191	0013	0020
740	.0363	0:154	.0349	0175	0014	0020
750	.0354	0137	•0340	0159	0014	0021



TABLE I.- SECTION COORDINATES - Concluded

x/c	(y/c) _u	(y/c) ₁	(y/c) _u	(y/c) ₁		1//
x/c	Airfo	il 31	Airfo	il 33	Δ(y/c) _u	Δ(y/c) _l
0.760	0.0345	-0.0121	0.0330	-0.0143	-0.0015	-0.0022
.770	.0336	0105	.0320	0128	0016	0023
.780	.0326	0089	.0309	0113	0017	0024
•790	.0315	0074	.0298	0099	0017	0025
.800	.0304	0060	.0287	0085	0017	0025
.810	.0292	0047	.0275	0072	0017	0025
.820	.0280	0035	.0262	0060	0018	0025
.830	.0267	0024	.0248	0049	0019	0025
.840	.0254	0014	.0234	0038	0020	0024
.850	.0240	0006	.0219	0029	0021	0023
.860	.0225	0.0	.0204	0022	0021	0022
.870	.0210	.0005	.0188	0017	0022	0022
.880	.0194	.0007	.0171	0014	0023	0021
.890	.0176	.0007	.0153	0013	0023	0020
•900	.0157	.0005	.0135	0013	0022	0018
.910	.0137	.0001	.0116	0016	0021	0017
•920	.0116	0005	.0096	0021	0020	0016
•930	.0093	0014	.0075	0028	0018	0014
•940	.0069	0026	.0054	0039	0015	0013
•950	.0044	0041	.0032	0053	0012	0012
•960	.0019	0059	.0008	0069	-,0011	0010
970	0008	0080	0017	0088	0009	0008
.980	0037	0105	0044	0110	0007	~.0005
.990	0068	0133	0074	0135	0006	0002
1.000		0164		0163		.0001
	1			•		

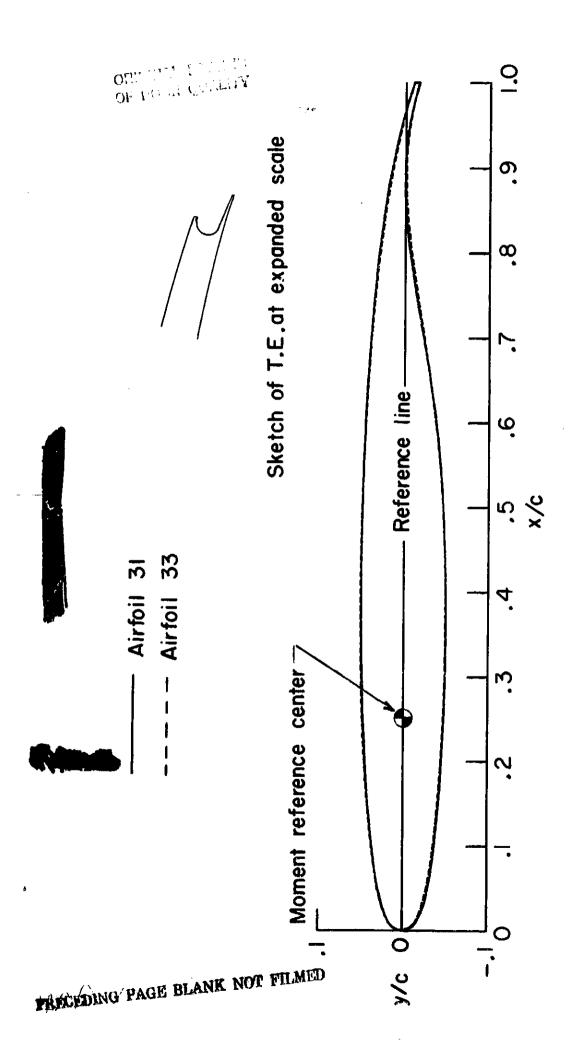


Figure I. - Sketches of 10-percent-thick supercritical airfoils 31 and 33.

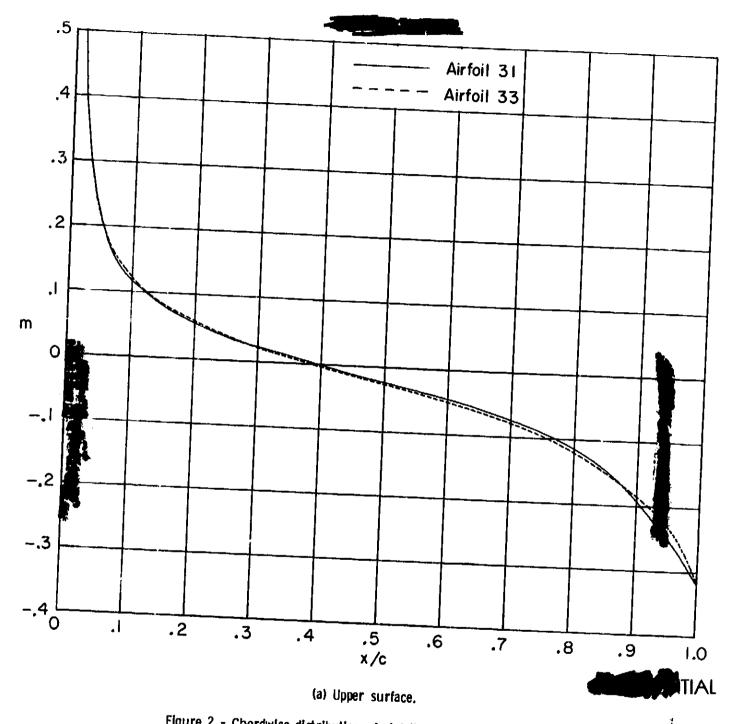


Figure 2. - Chordwise distribution of airfoil surface slopes and curvature.

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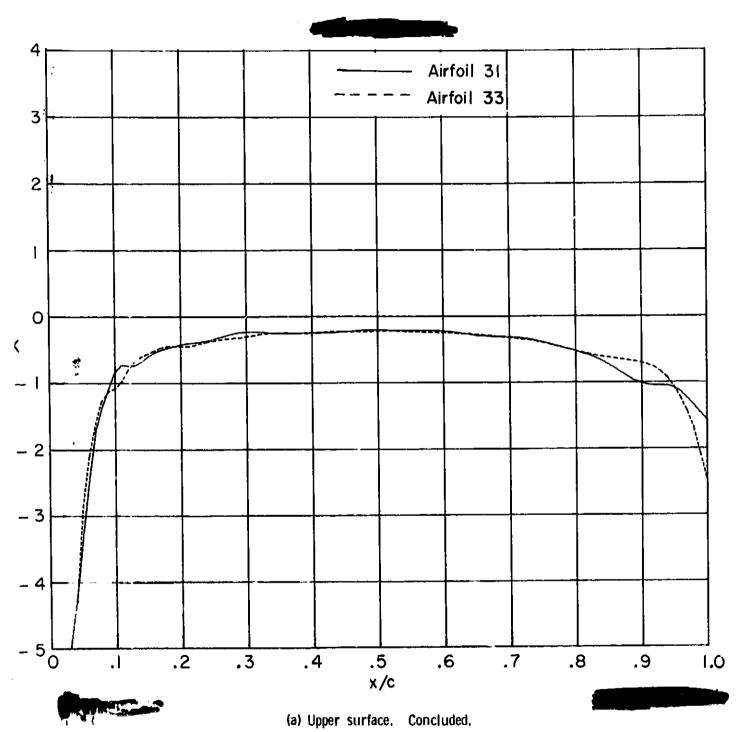


Figure 2. - Continued.

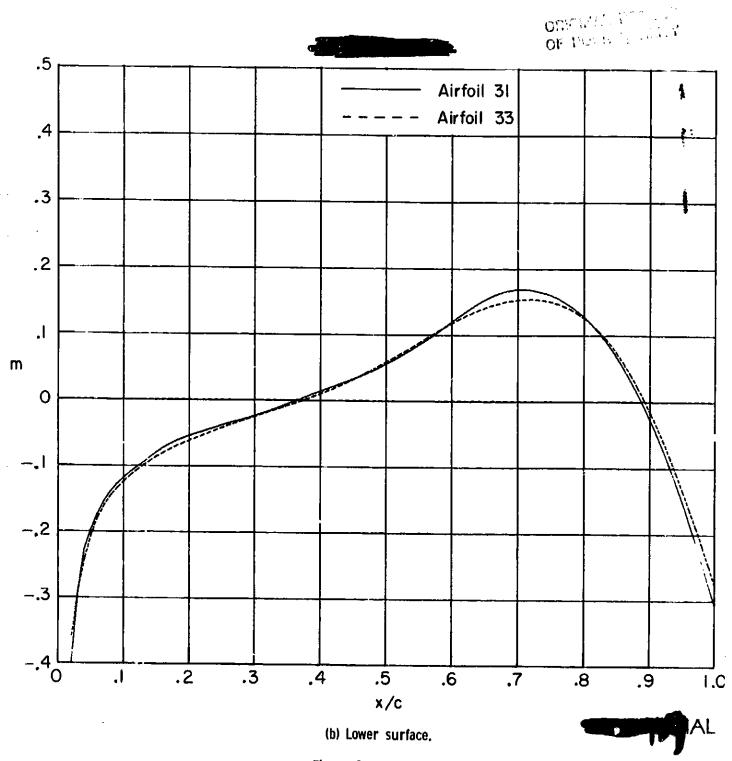


Figure 2. - Continued.

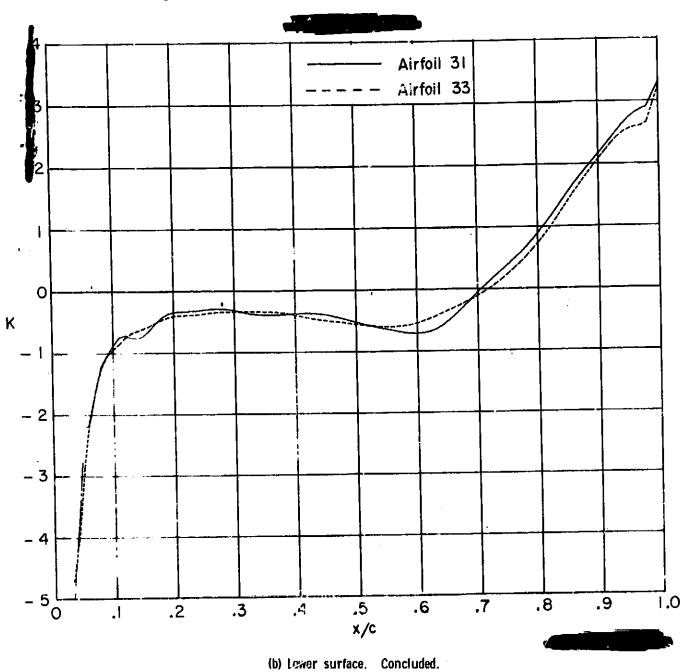
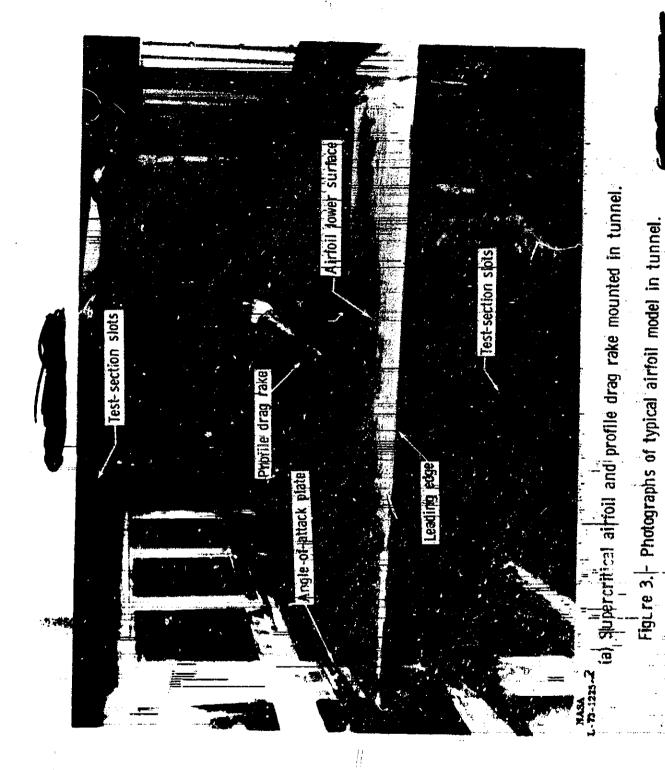
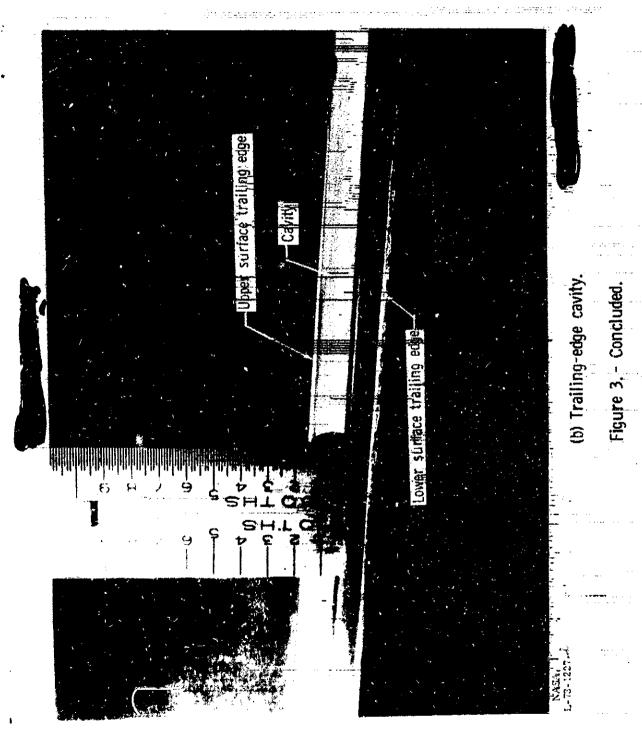


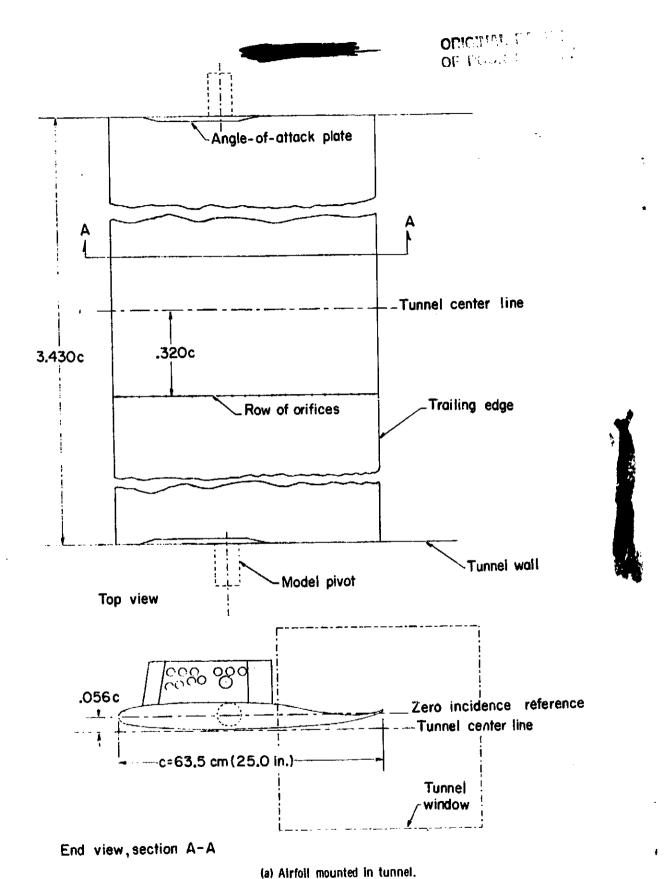
Figure 2. - Concluded.



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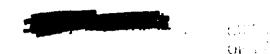


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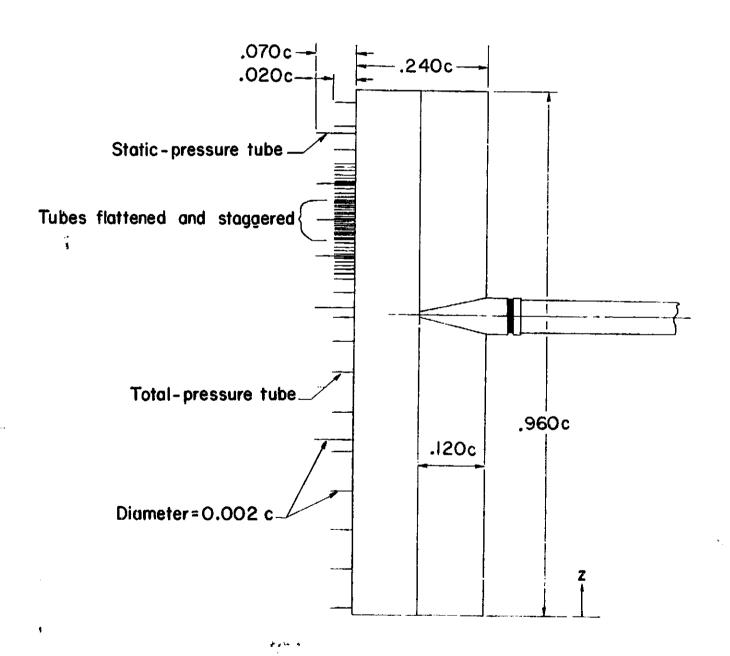


(a) Allion induited in termo.

Figure 4. - Apparatus dimensions in terms of chord. c = 63.5 cm (25.0 in.).







(b) Profile drag rake.

Figure 4. - Concluded.



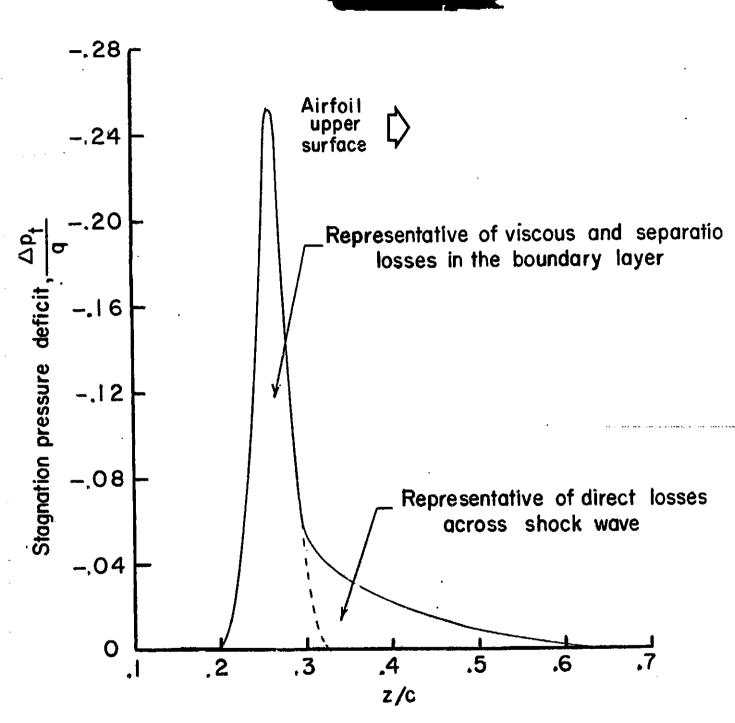


Figure 5. - Schematic of wake profiles.

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8 x 10⁶
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5.50 .60 .70 .80 .90

Figure 6. - Variation with Mach number of test wind-tunnel Reynolds number.

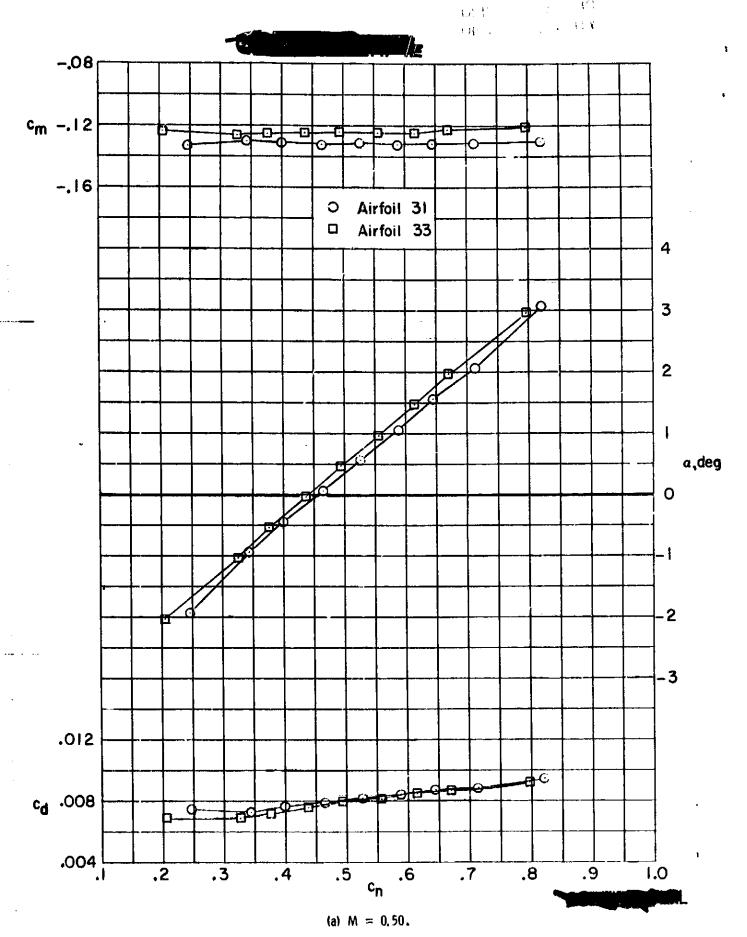
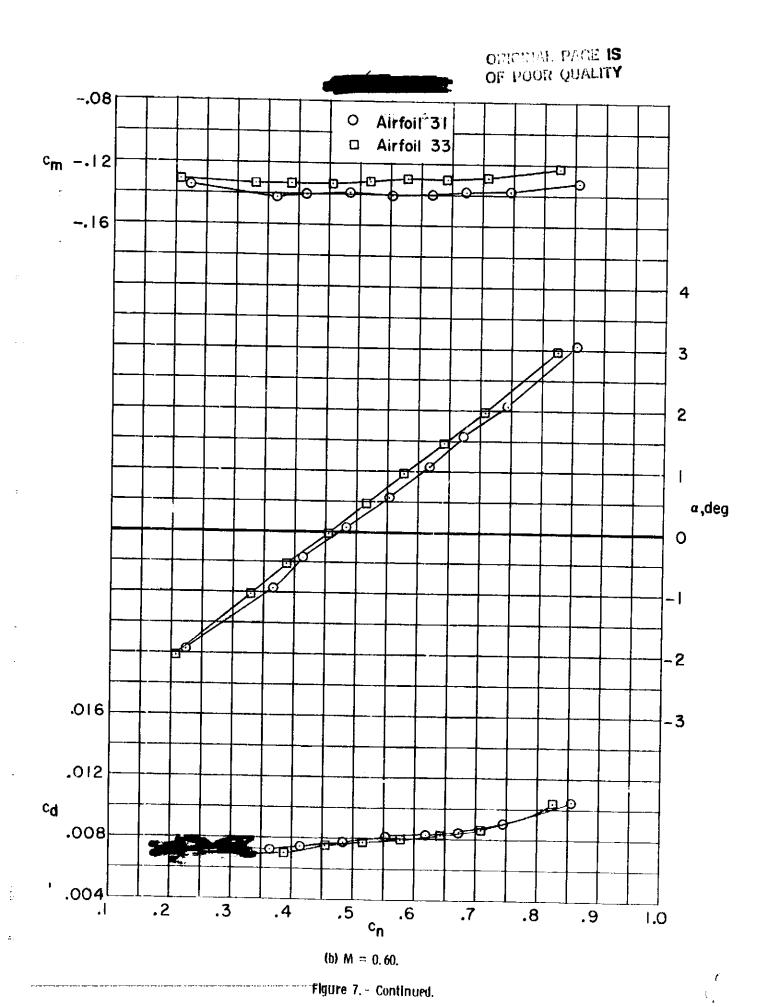
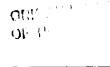


Figure 7. - Comparison of force and moment characteristics of 10-percent-thick supercritical airfolis 31 and 33.





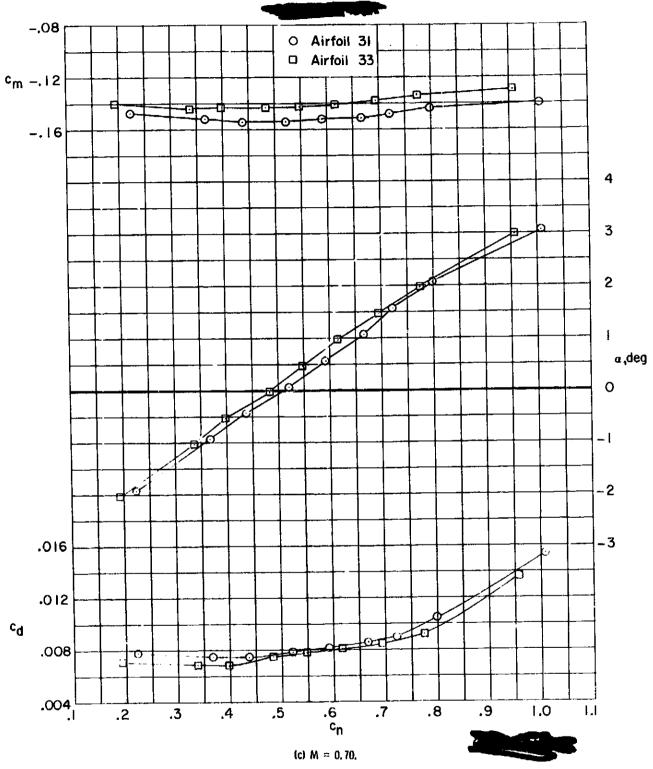
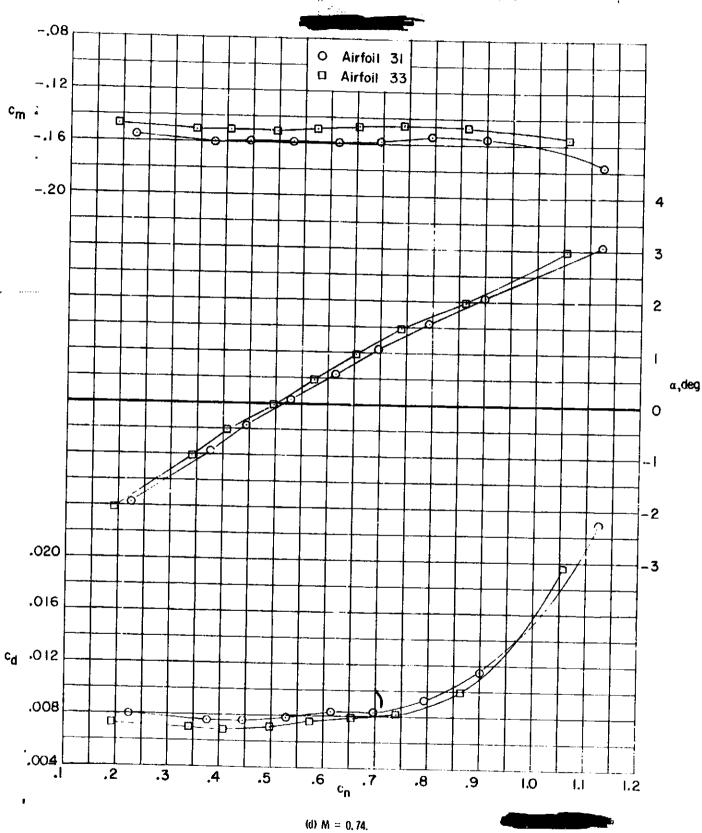


Figure 7. - Continued.



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Figure 7. - Continued.

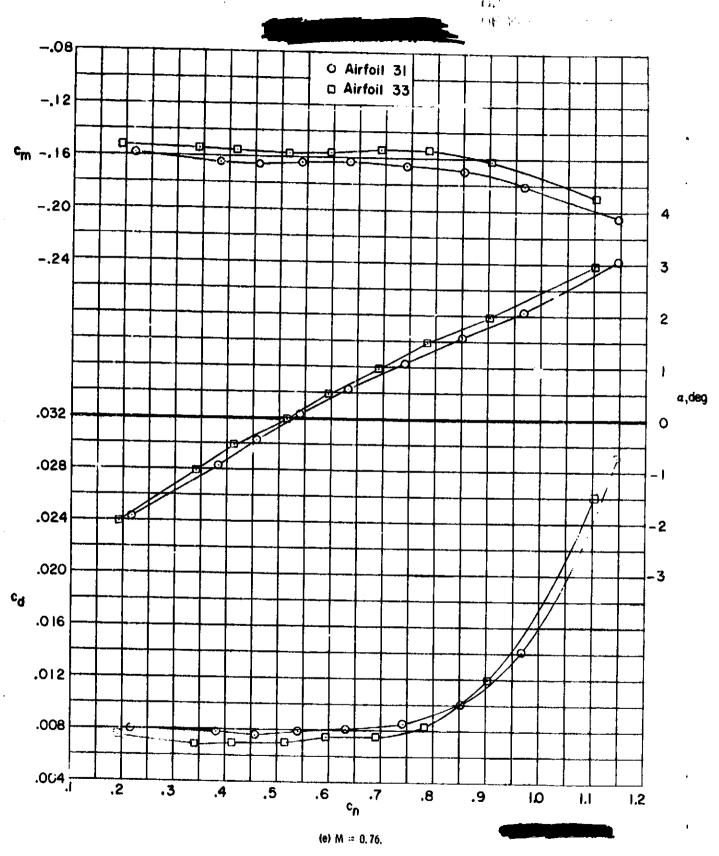


Figure 7. - Continued.

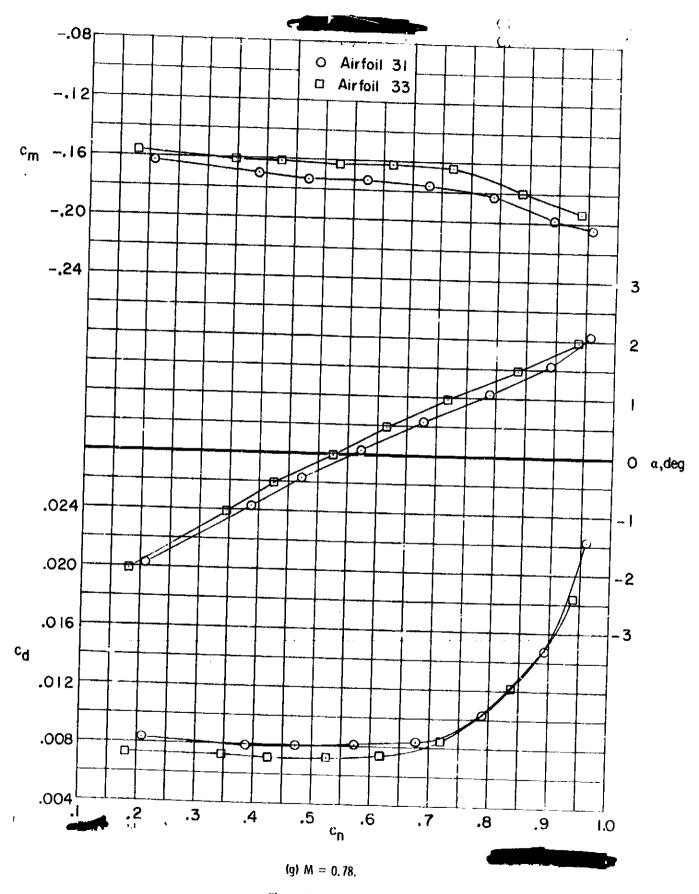


Figure 7. - Continued.

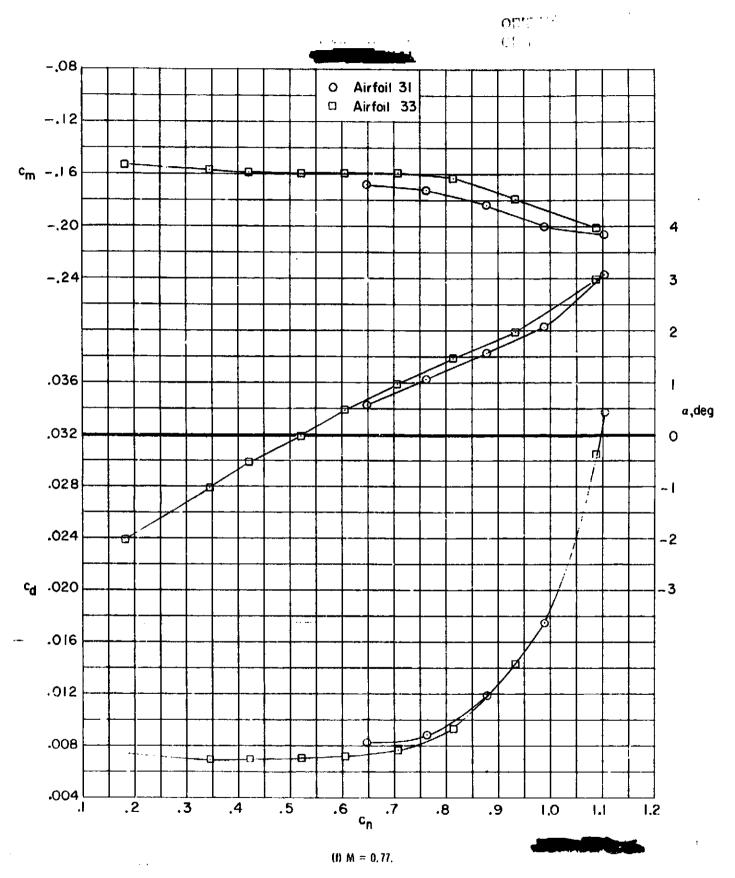


Figure 7. - Continued.

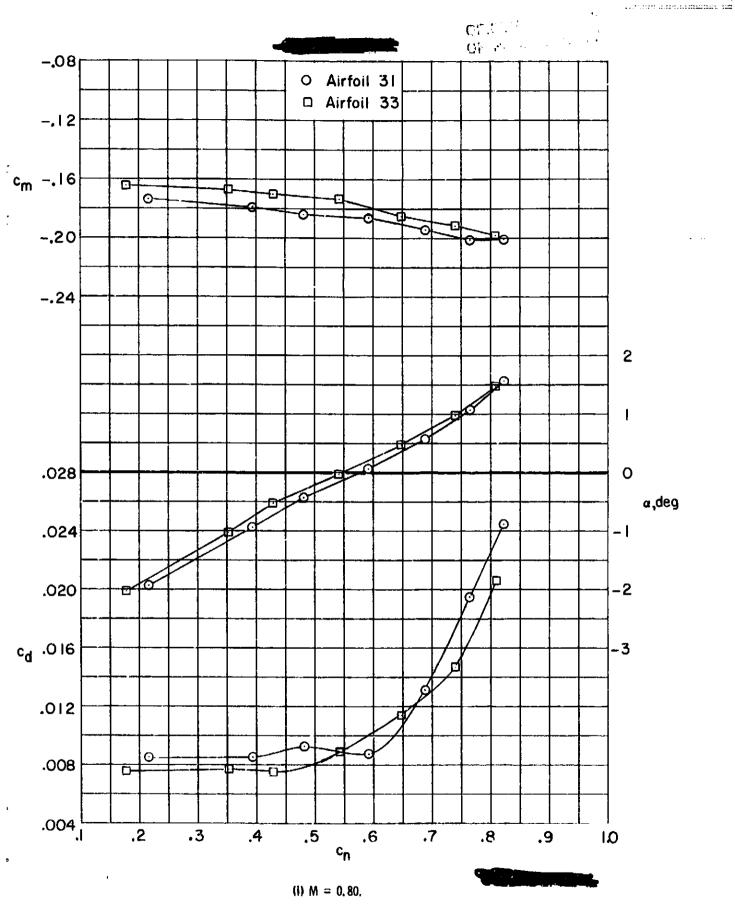


Figure 7. - Continued.

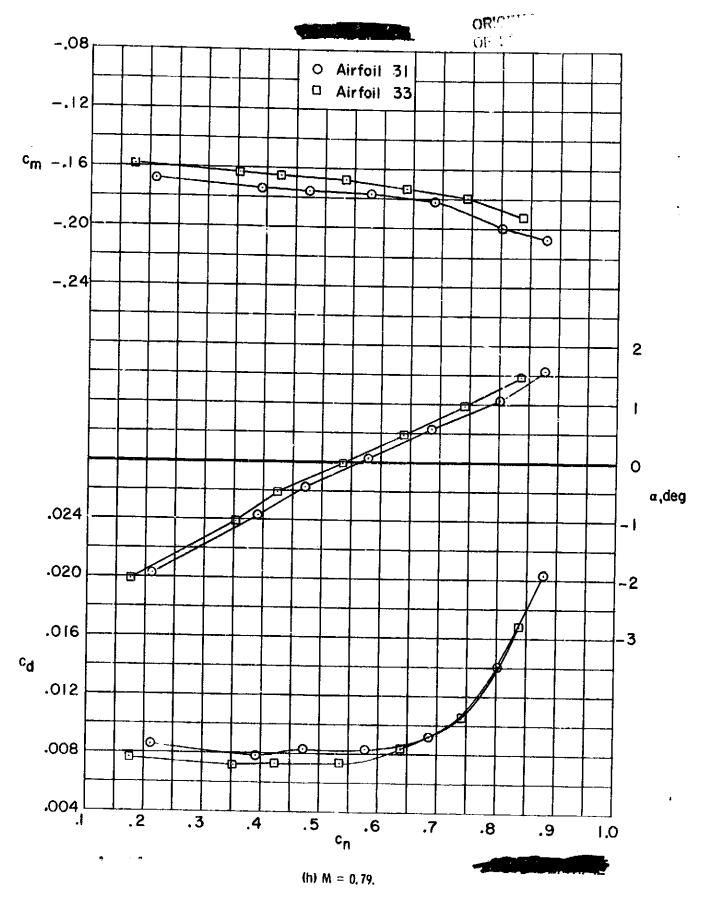


Figure 7. - Continued.

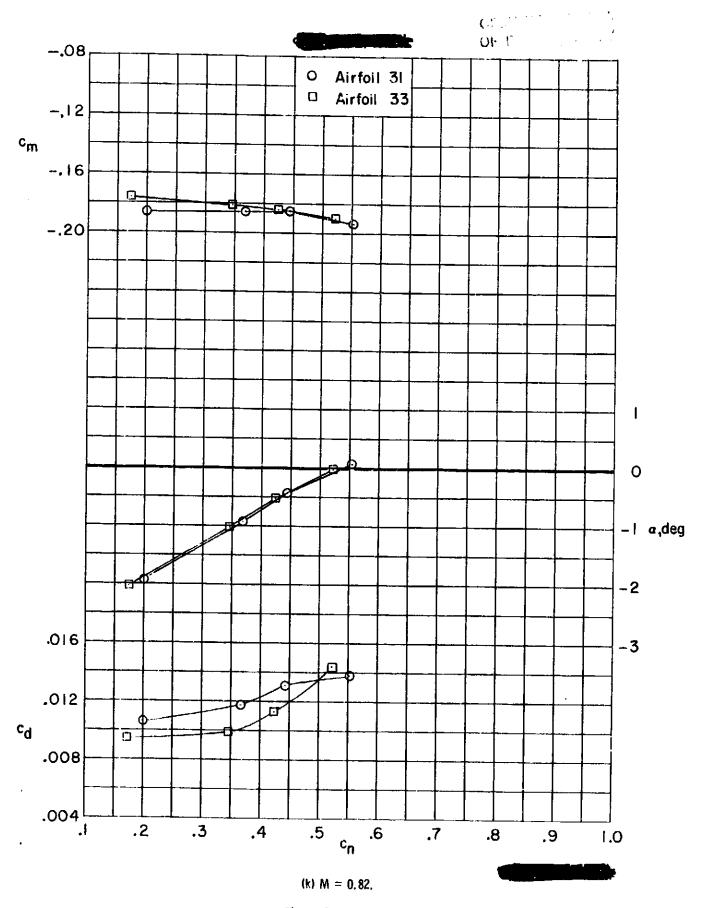


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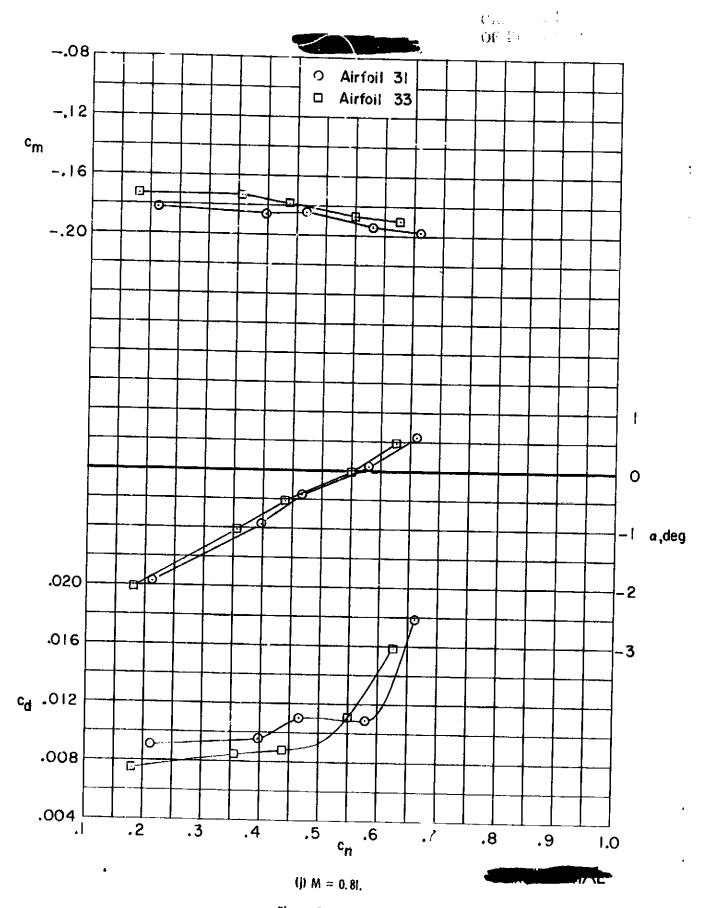


Figure 7. - Continued.

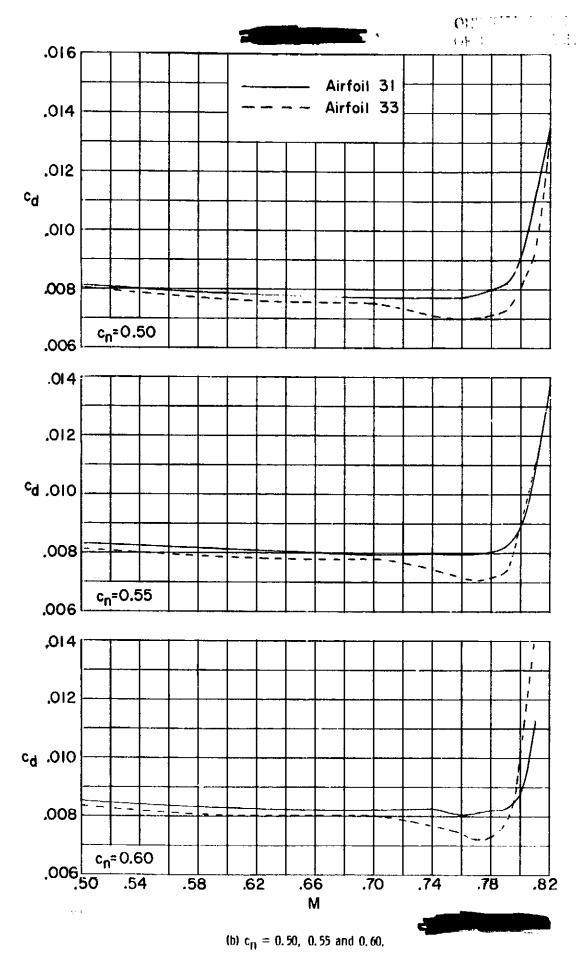


Figure 8. - Continued,

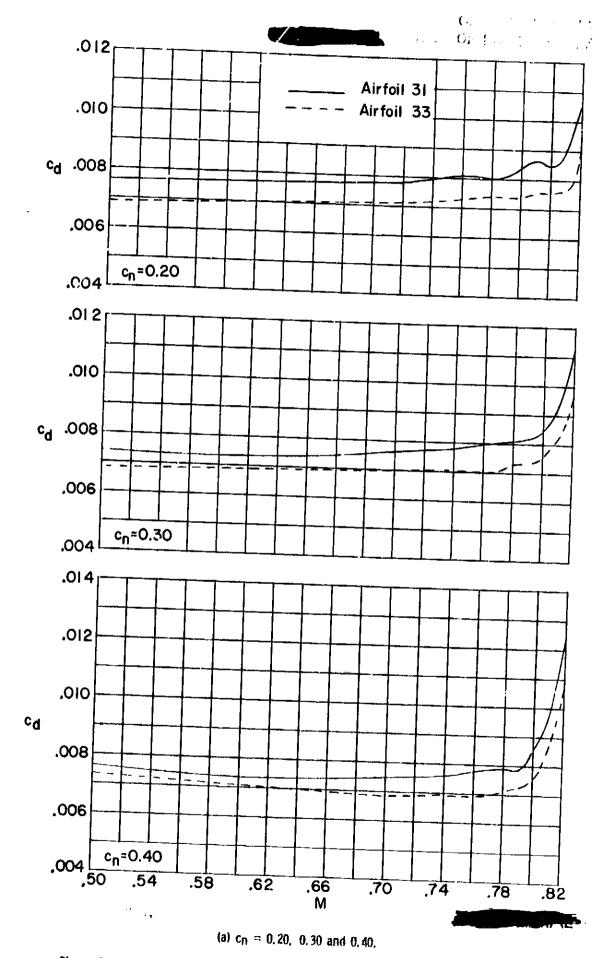
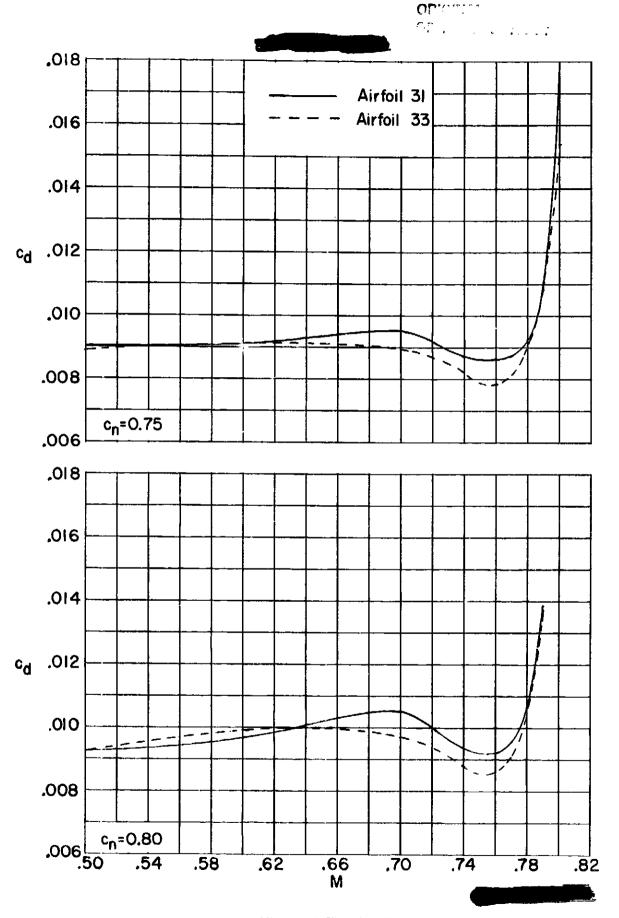


Figure 8. - Variation of measured section drag coefficient with Mach number of 10-percent-thick supercritical airfolls 31 and 33.



(d) $c_n = 0.75$ and 0.80.

Figure 8. - Continued.

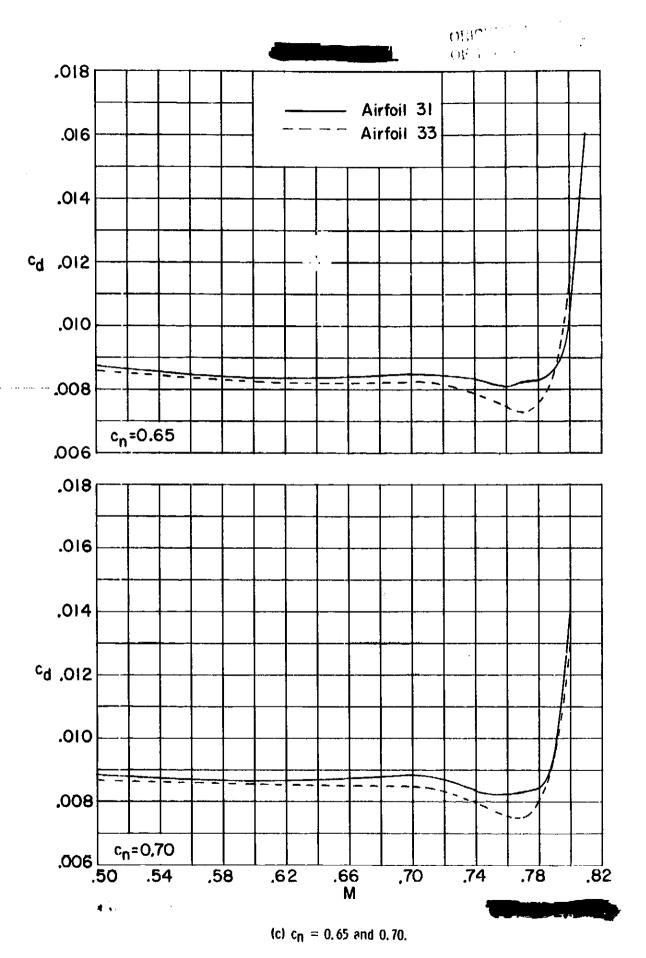


Figure 8. - Continued.

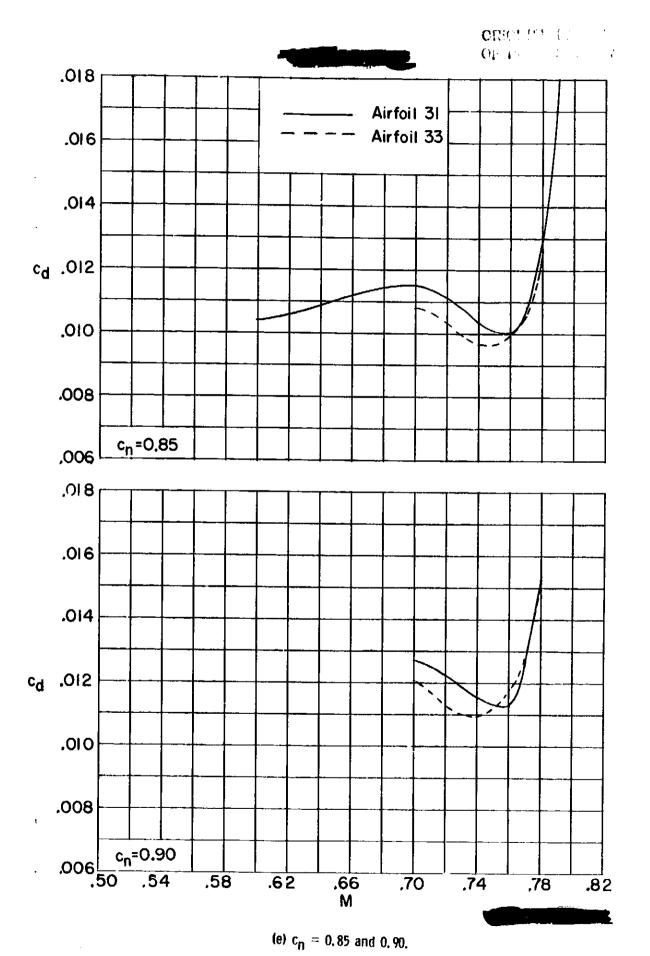


Figure 8. - Concluded.

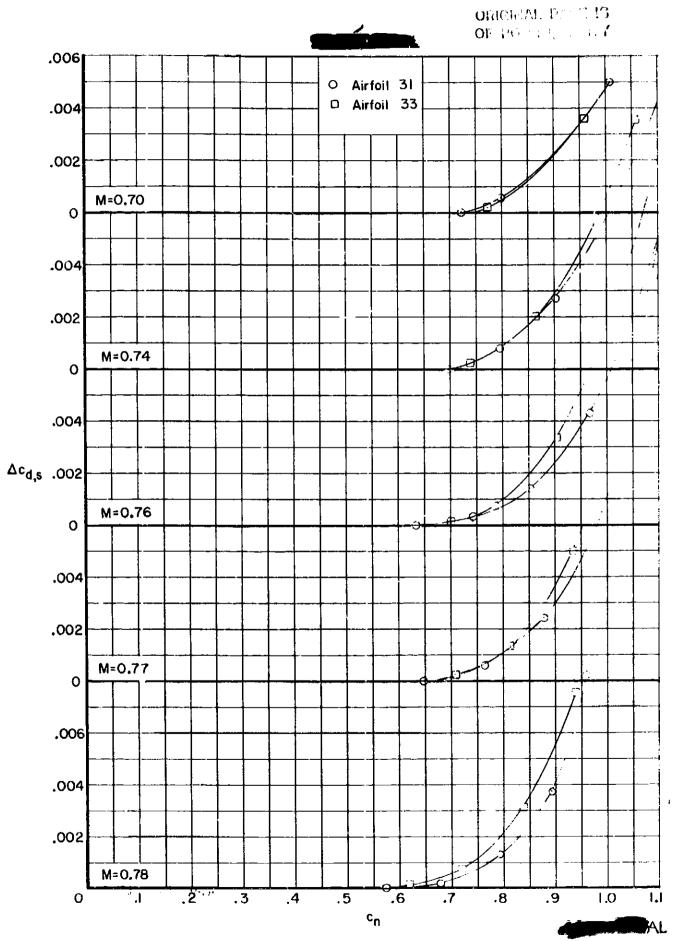


Figure 9. - Drag increment due to shock wave losses.

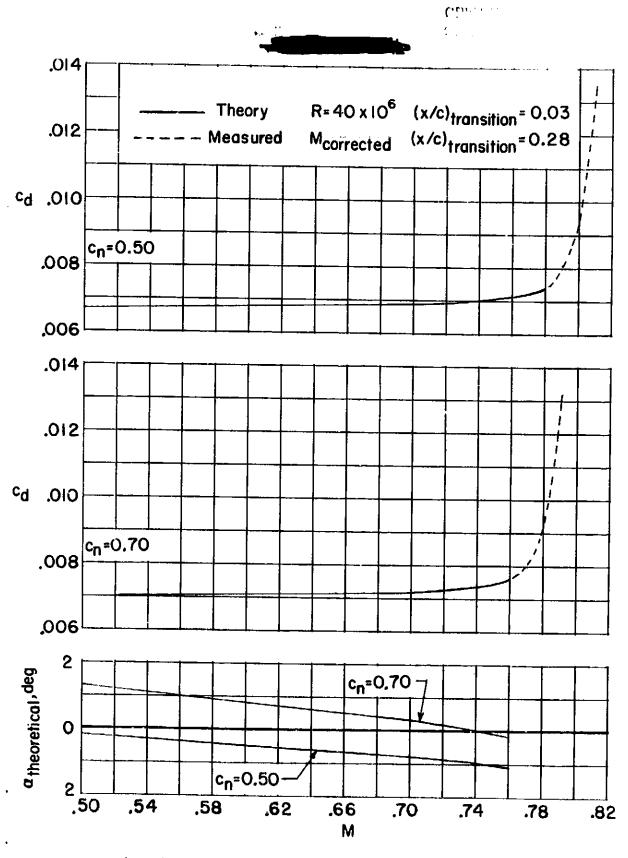


Figure 10. - Composite drag-rise curves and theoretical angles of attack for 10-percent-thick supercritical airfoil 33.

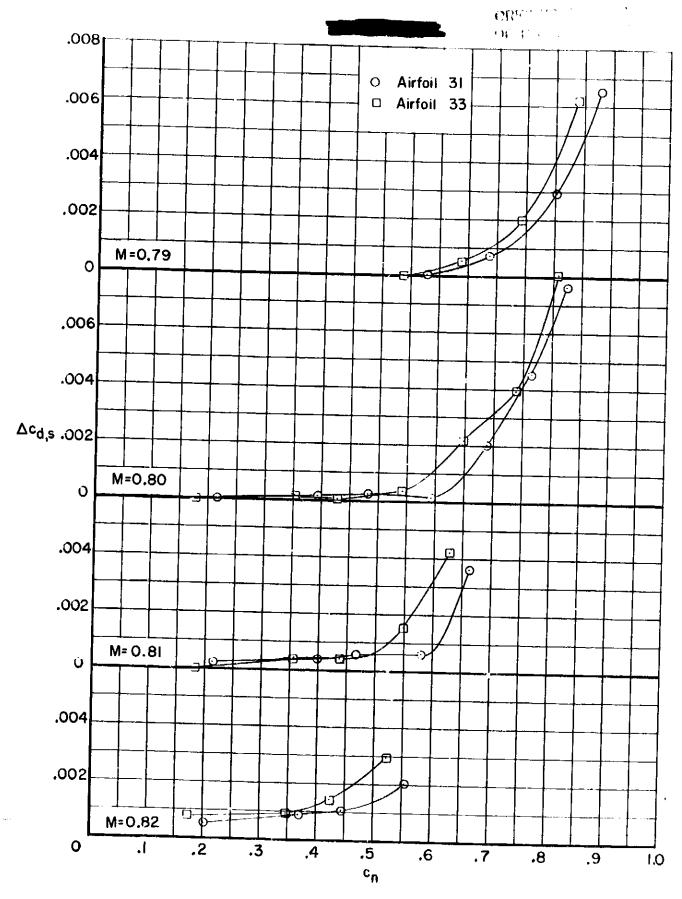
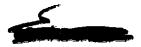
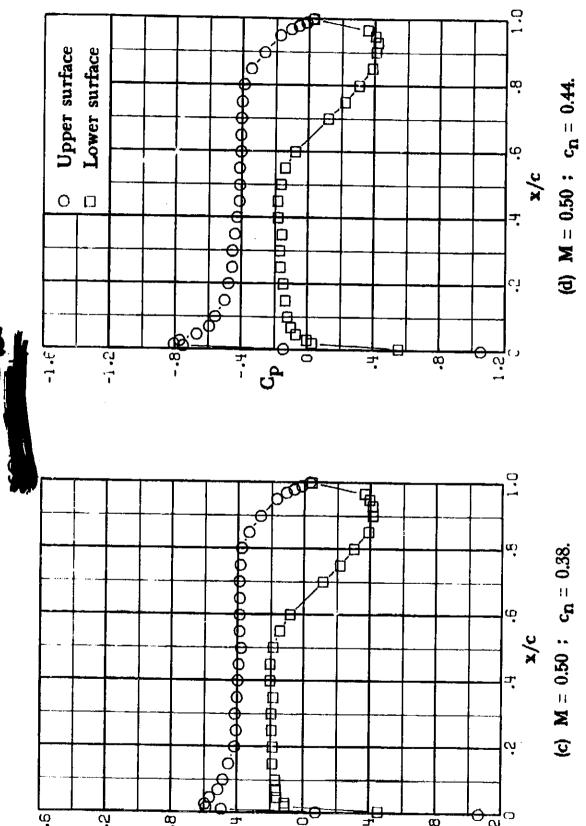


Figure 9. - Concluded.





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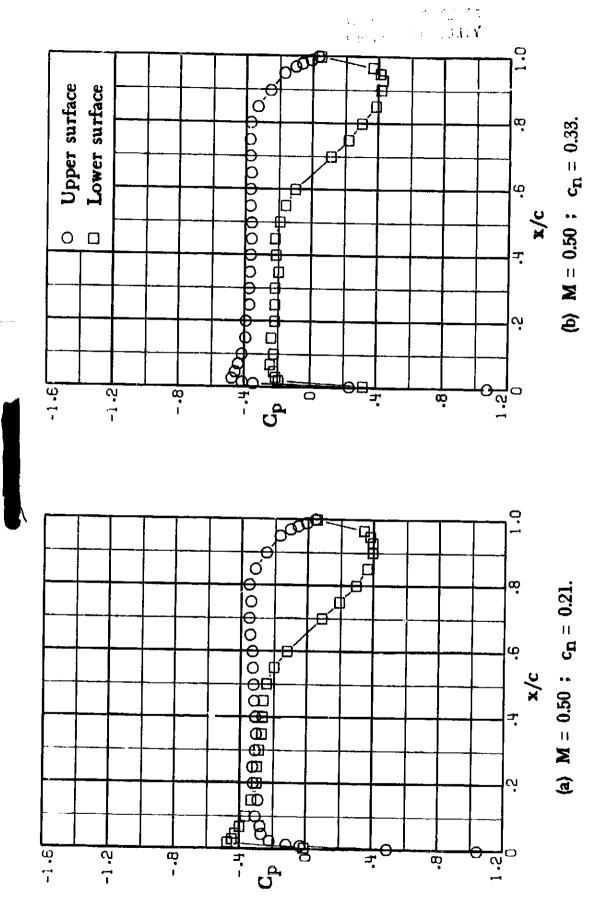
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Figure II. - Continued.

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M = 0.50

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Figure II. – Chordwise pressure

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Figure 11. - Continued.

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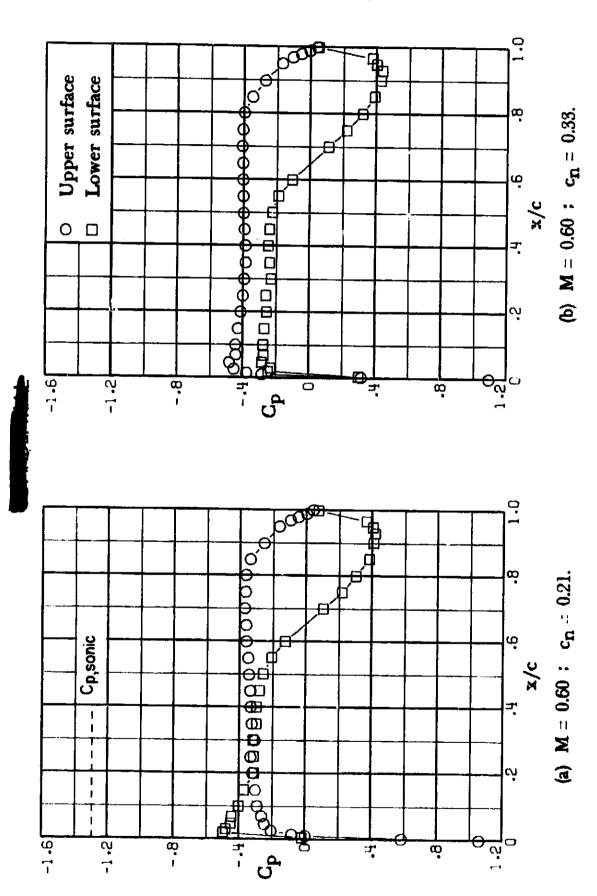
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Figure II. - Continued

(e) M = 0.50; $c_{m} = 0.49$.



M = 0.60Figure 12. - Chordwise pressure distributions for 10-percent thick supercritical airfoil 33.

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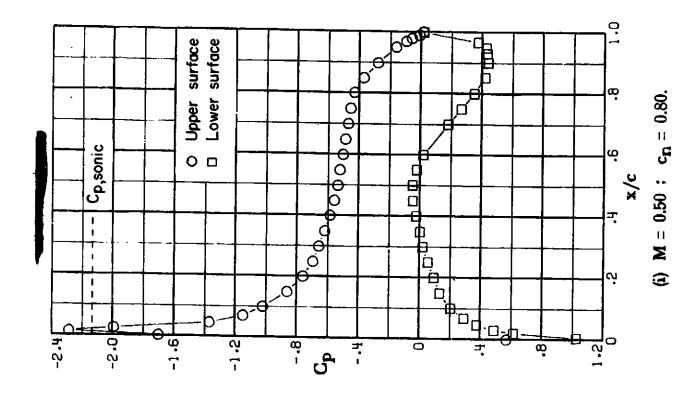


Figure II. - Concluded.

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Figure 12. - Continued.

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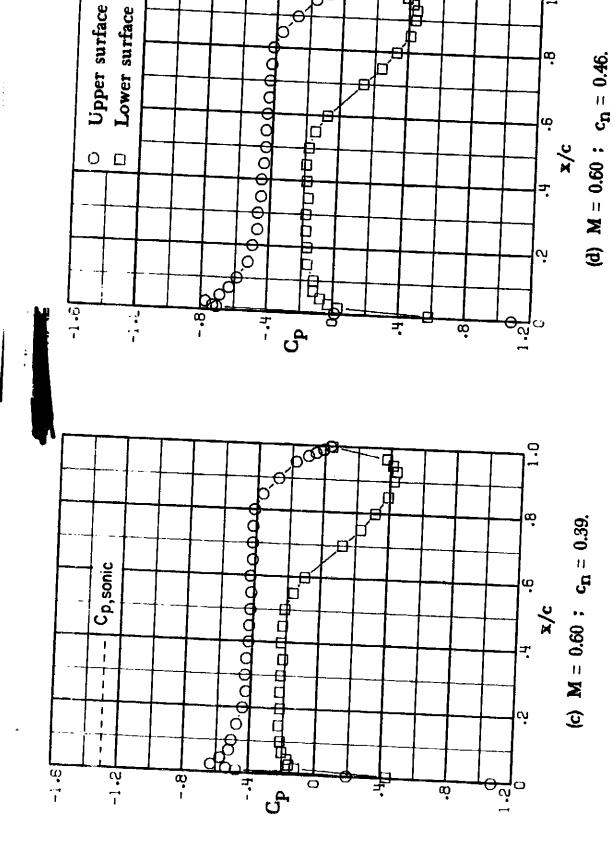


Figure 12. - Continued.

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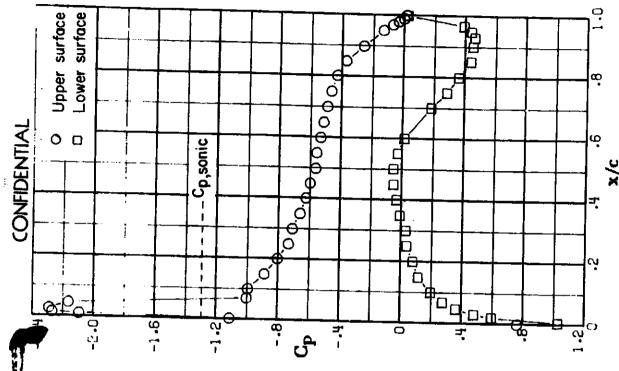
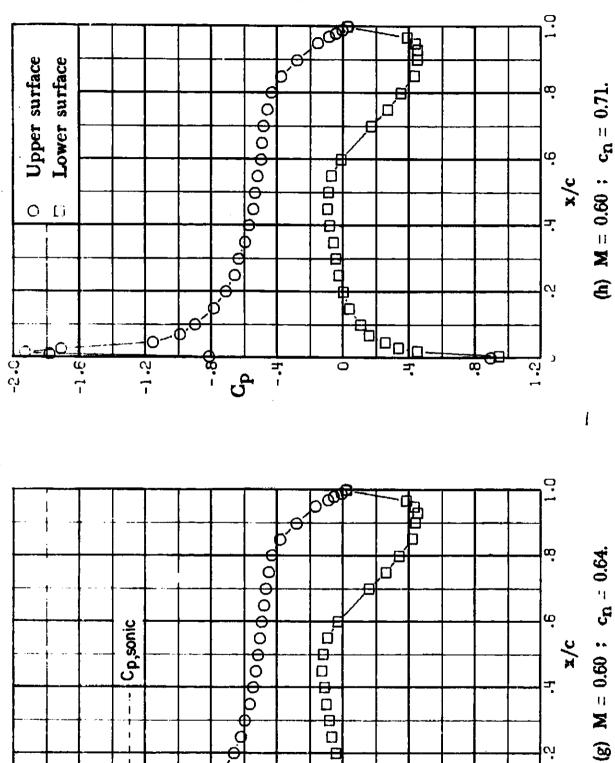


Figure 12. - Concluded.

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(i) M = 0.60; $c_n = 0.82$.



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Figure 12. - Continued

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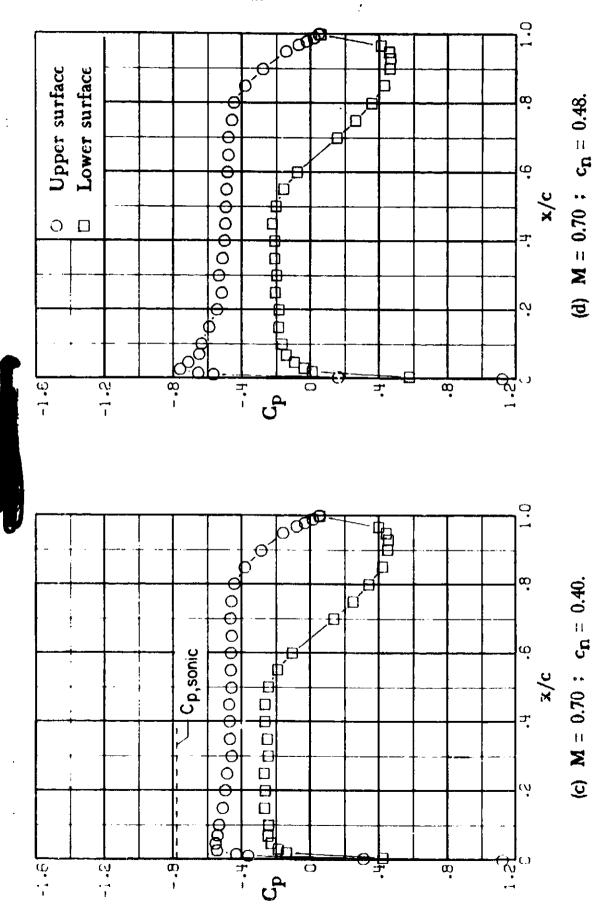
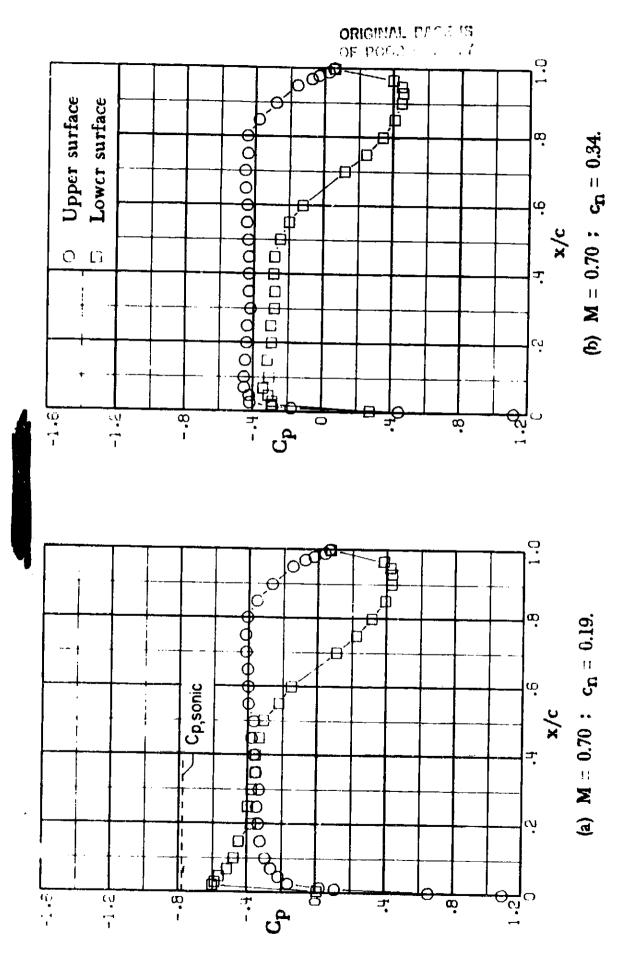


Figure 13. - Continued.



M = 0.70Figure 13. - Chordwise pressure distributions for 10-percent thick supercritical airfoil 33.

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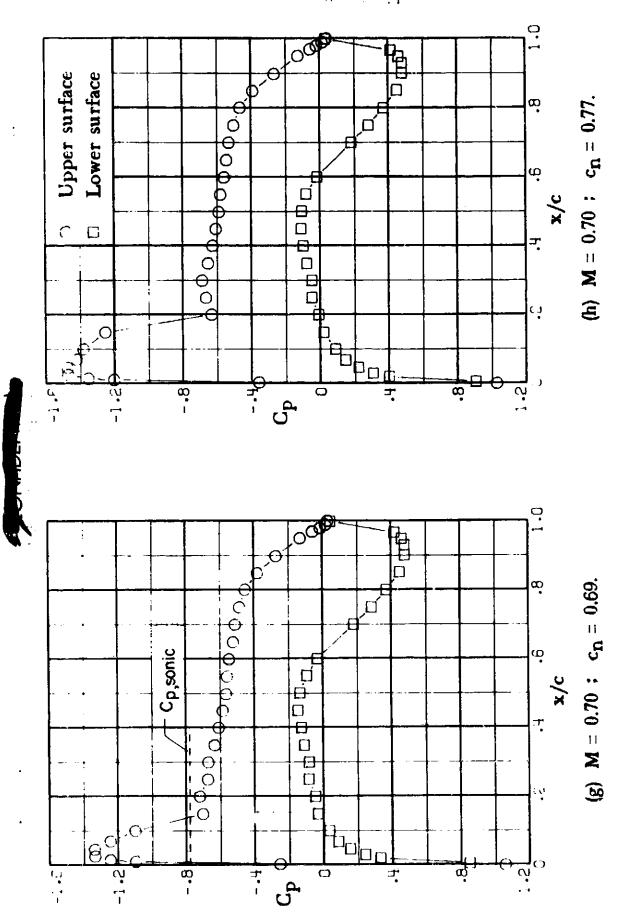


Figure 13. - Continued.

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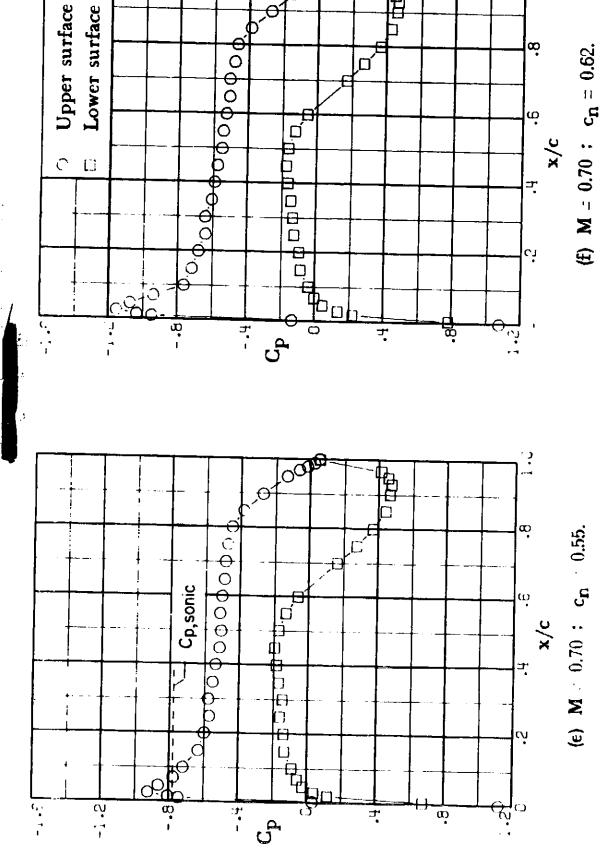
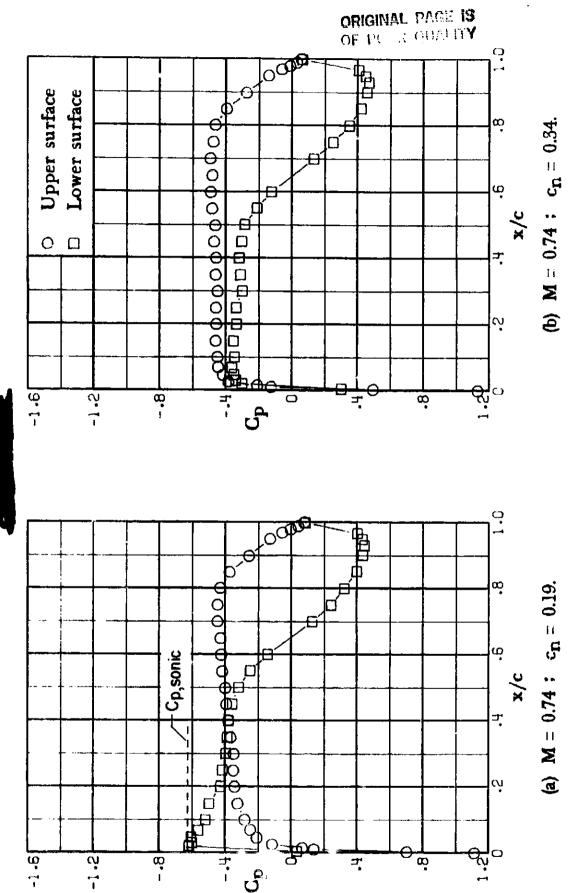


Figure 13. - Continued.





M = 0.74Figure 14. - Chordwise pressure distributions for 10-percent thick supercritical airfoil 33.

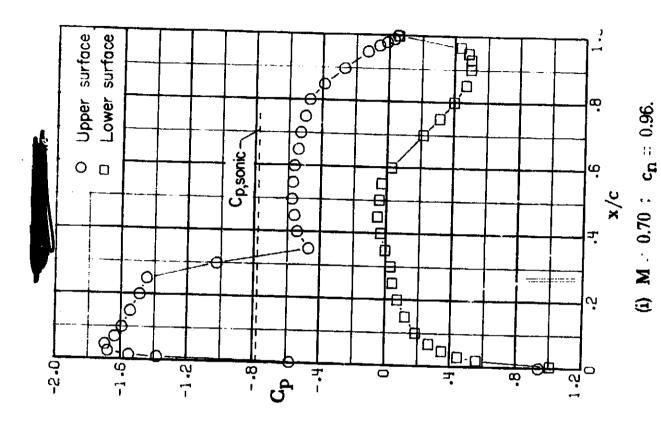


Figure 13. - Concluded.

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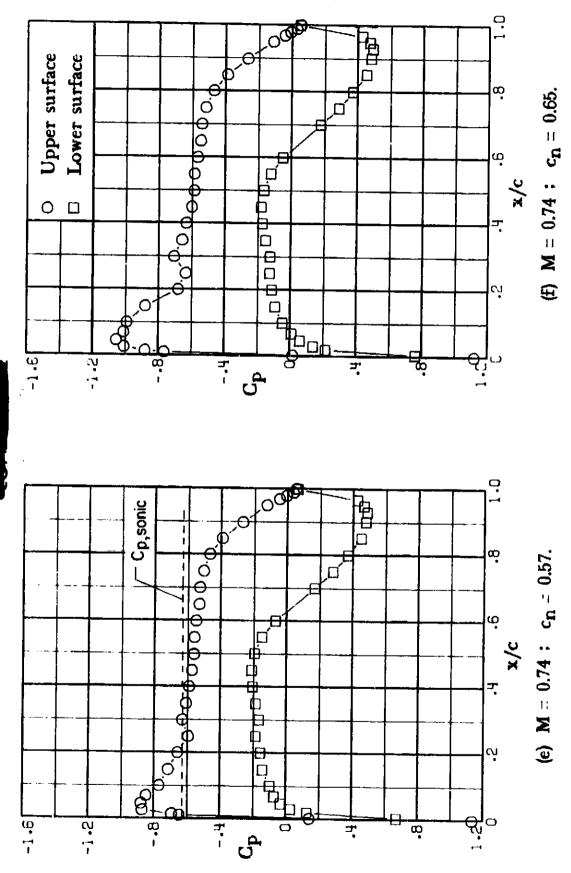


Figure 14. - Continued.

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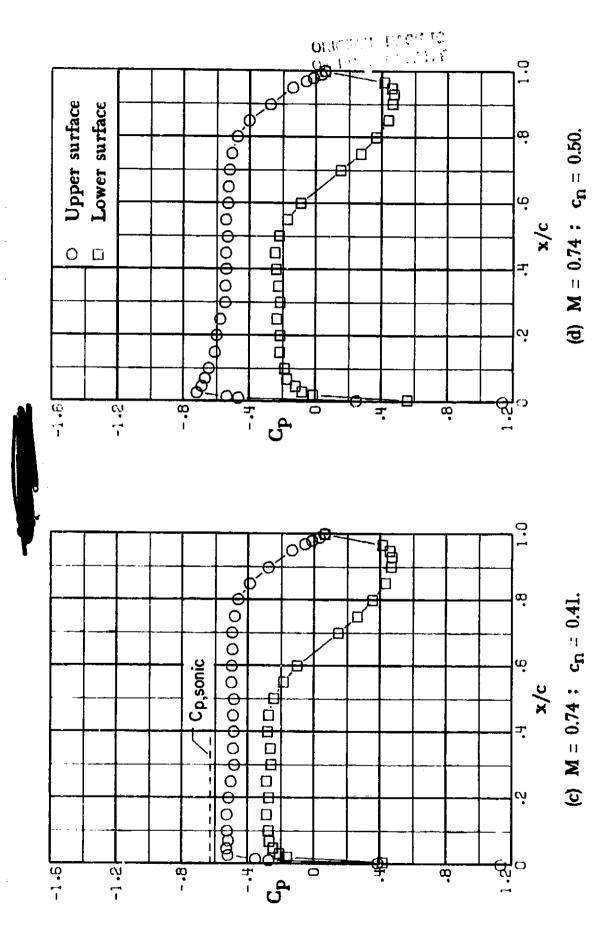


Figure 14. - Continued,

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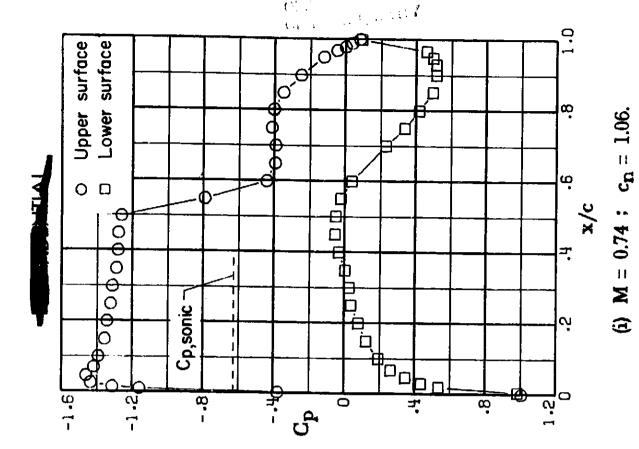
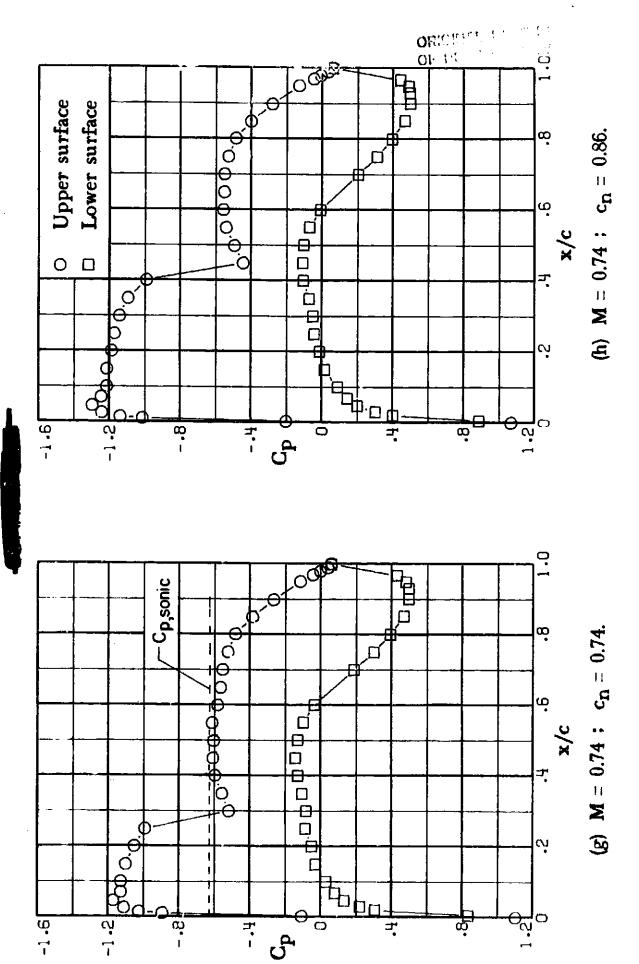


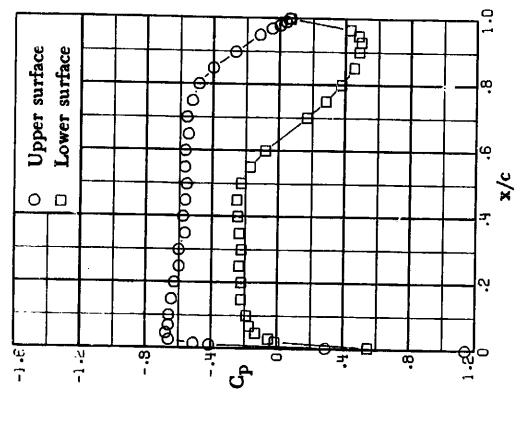
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Figure 14. - Continued.



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(d) M = 0.76; $c_n = 0.52$.

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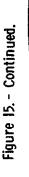
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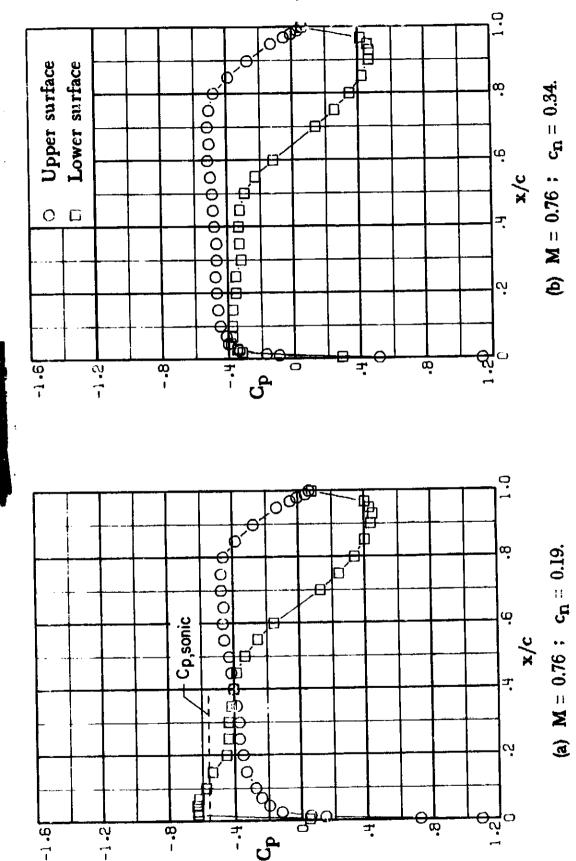
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(c) M = 0.76; $c_n = 0.41$.



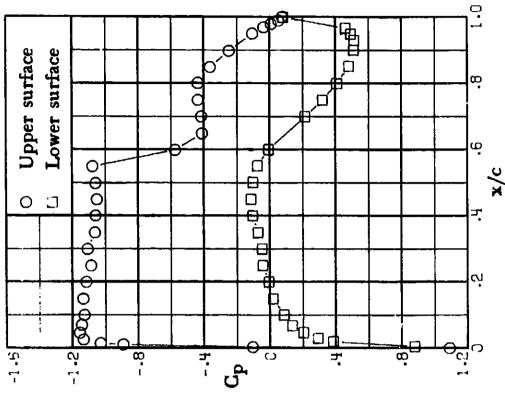
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M = 0.76Figure 15. - Chordwise pressure distributions for 10-percent thick supercritical airfoil 33.



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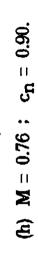


Figure 15. - Continued.

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(g) M = 0.76: $c_n = 0.78$.

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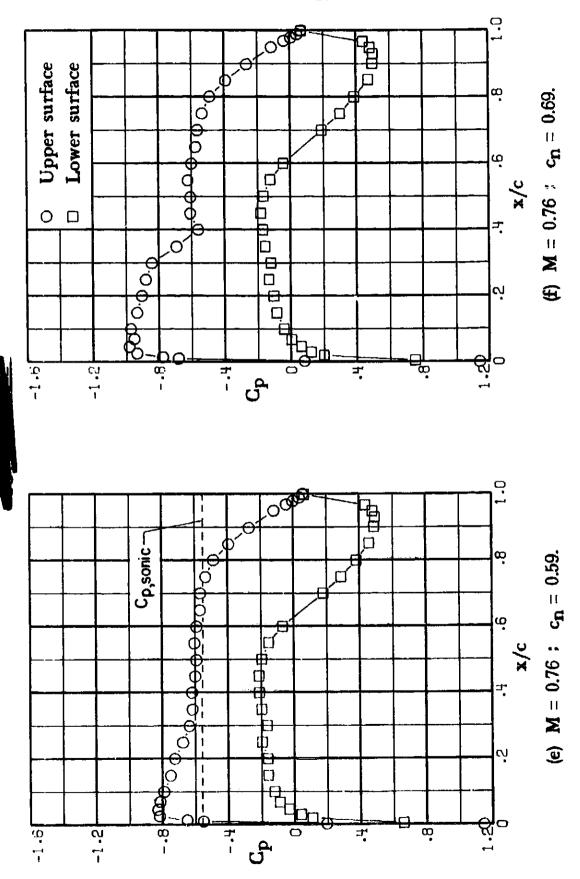


Figure 15. - Continued.



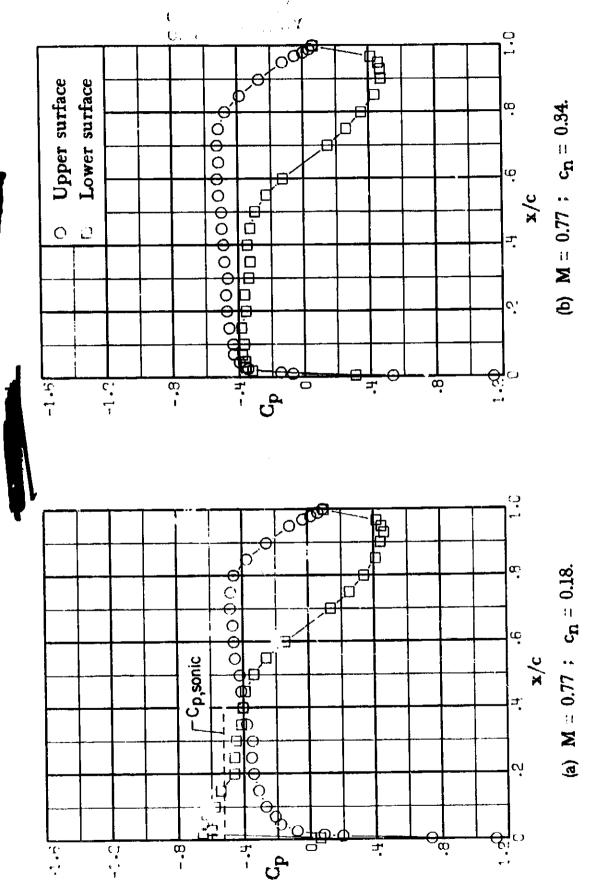


Figure 16. - Chordwise pressure distributions for 10-percent thick supercritical airfoil 33. M=0.77.

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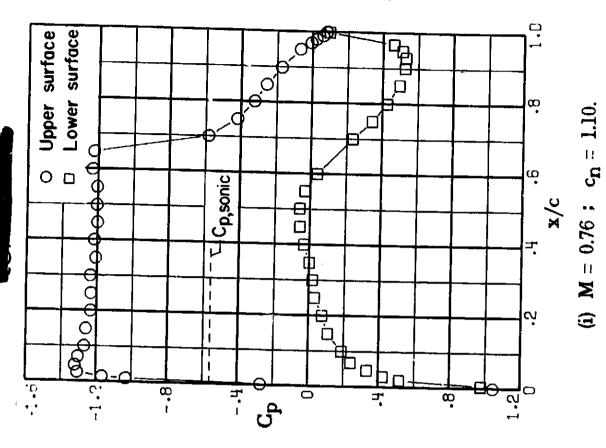


Figure 15. - Concluded.

Figure 15. - Concluded.

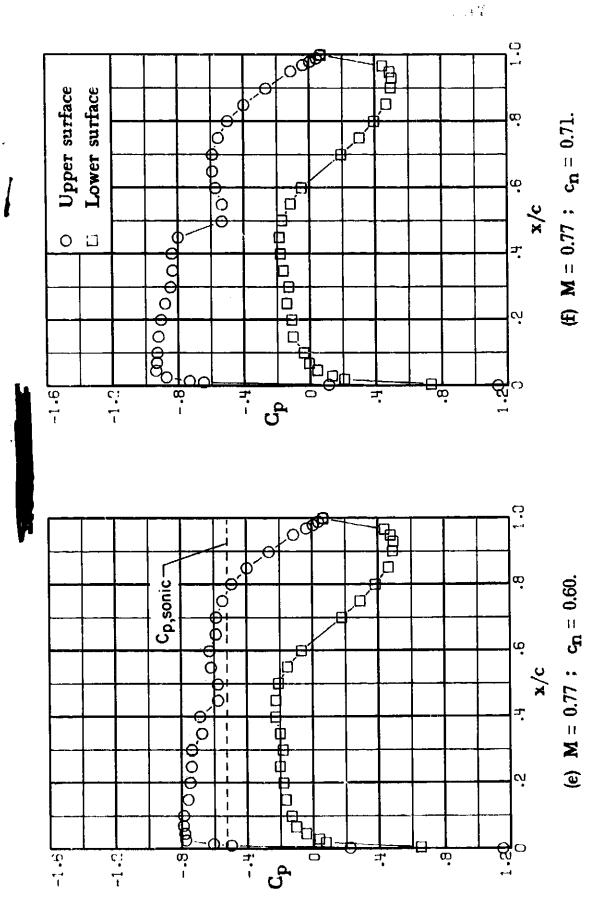
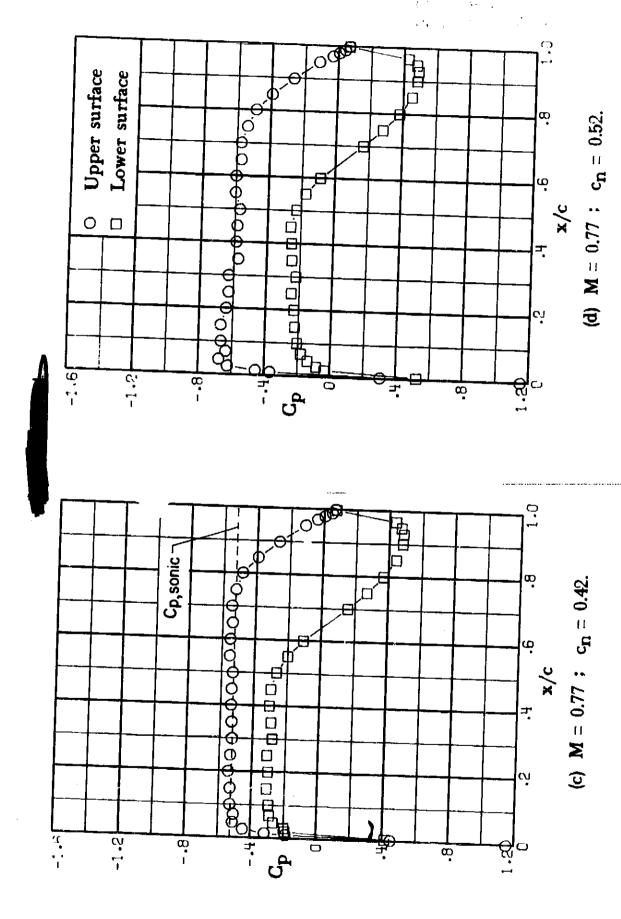


Figure 16. - Continued.

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Figure 16. - Continued.

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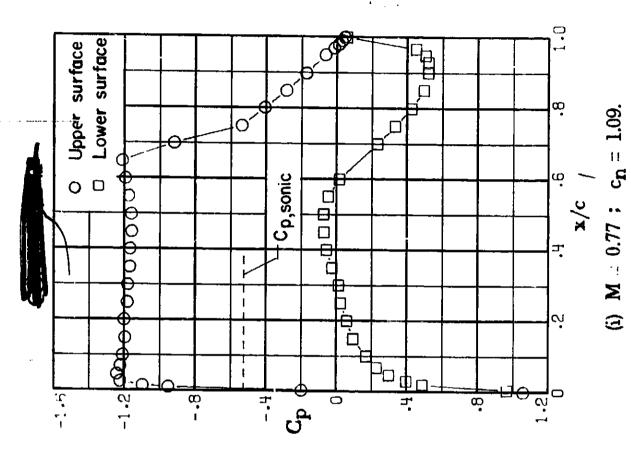


Figure 16. - Loncluded.

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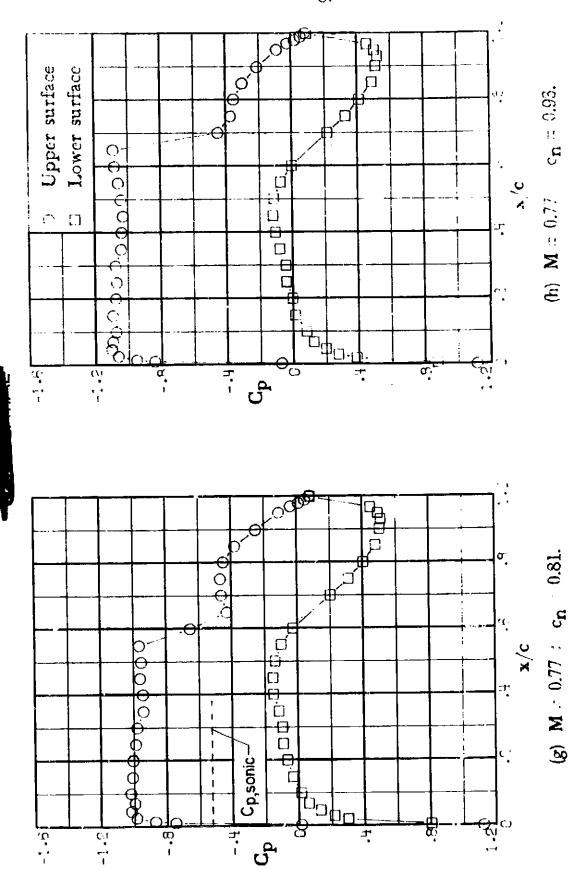
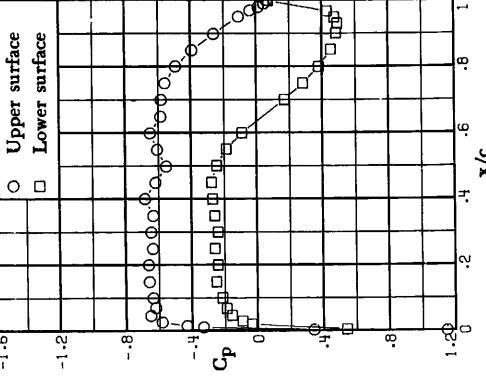


Figure 16. - Continued.



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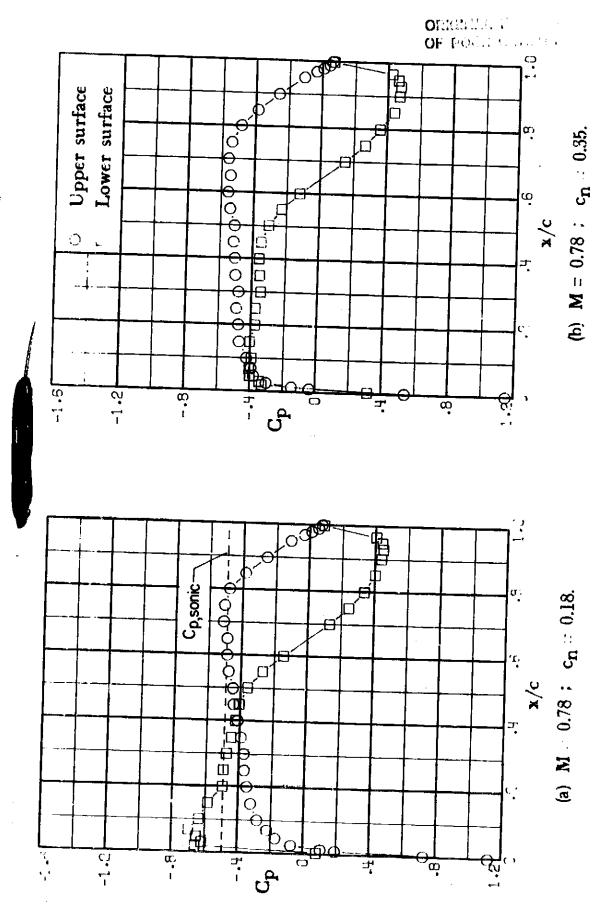
Figure 17. - Continued.

(d) M = 0.78; $c_n = 0.52$.

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(c) M = 0.78; $c_n = 0.42$.

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M = 0.78Figure 17.. - Chordwise pressure distributions for 10-percent thick supercritical airfoil 33.

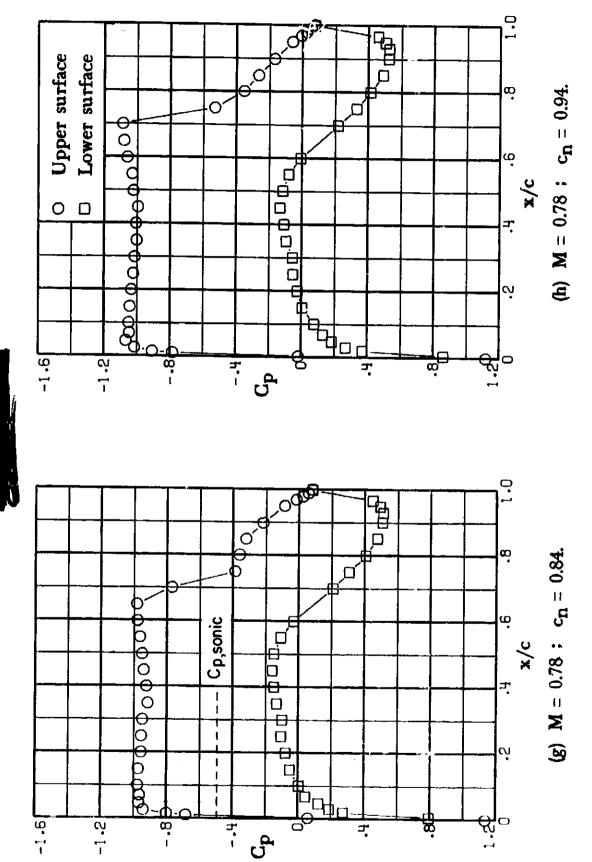
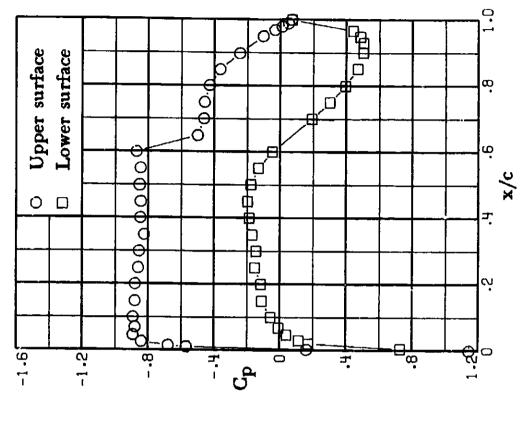


Figure 17. - Concluded.

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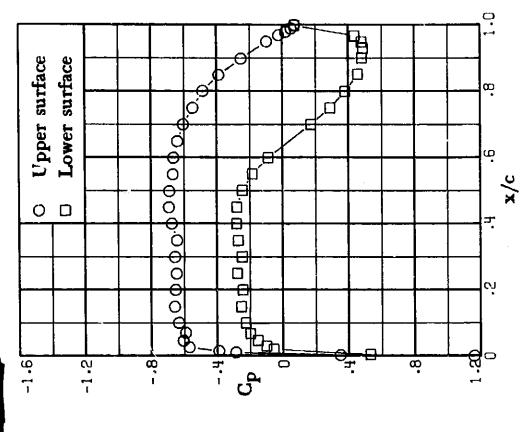
(e) M = 0.78; $c_n = 0.62$.

Figur: 17. - Continued.

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(d) M = 0.79; $c_n = 0.54$.

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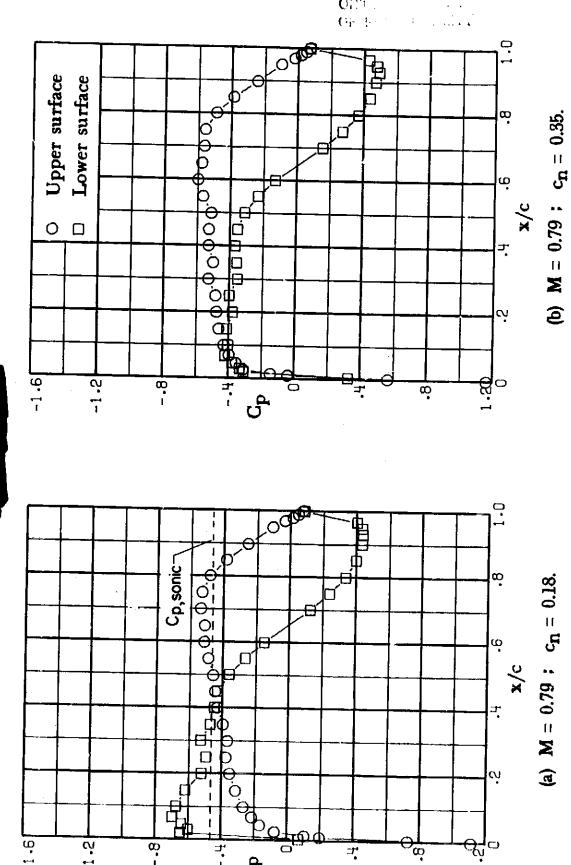
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(c) M = 0.79; $c_n = 0.42$.

Figure 18. - Continued.

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Figure 18. - Chordwise pressure distributions for 10-percent thick supercritical airfoll 33. M=0.79.

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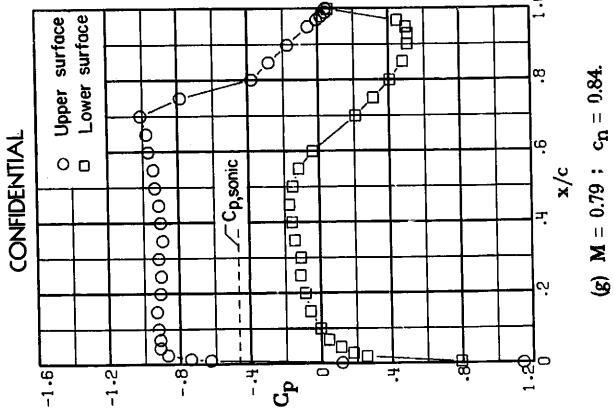


Figure 18. - Concluded.

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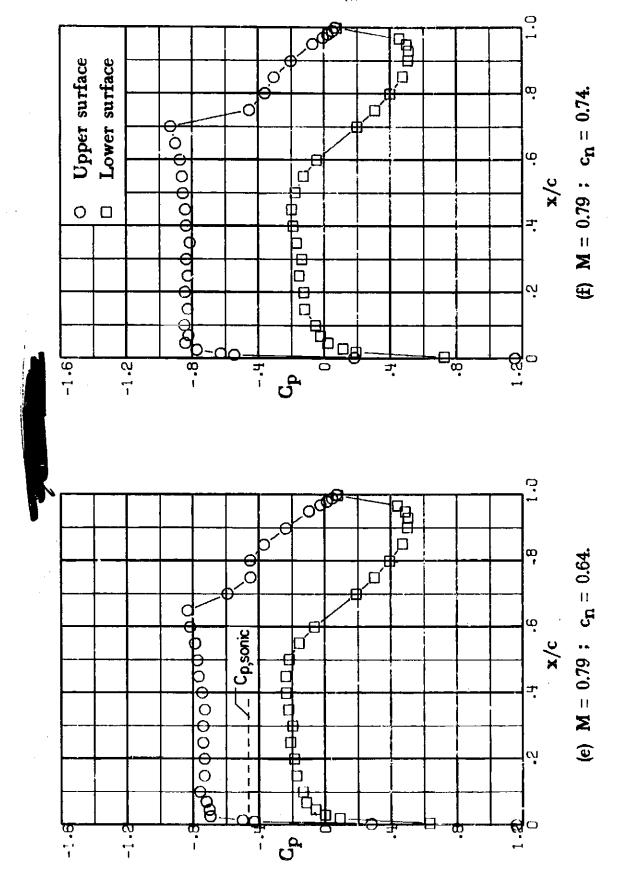


Figure 18. - Continued.

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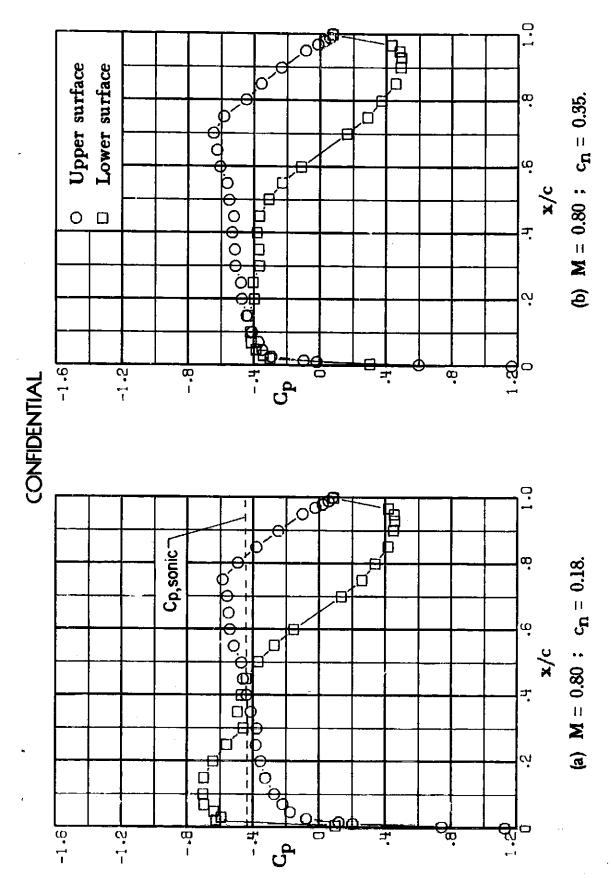


Figure 19. - Chordwise pressure distributions for 10-percent thick supercritical airfoil 33. M = 0.80.

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Figure 19. - Continued.

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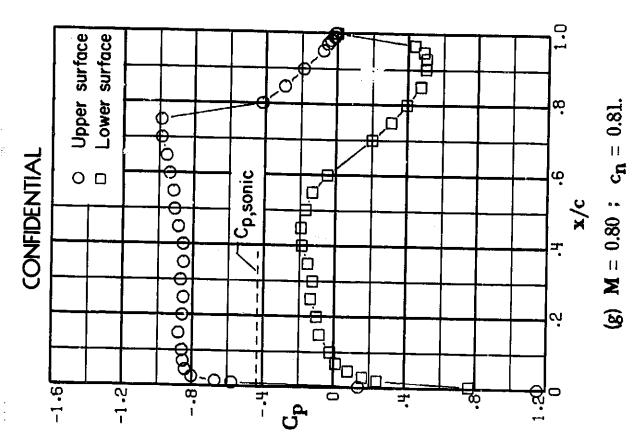


Figure 19. - Concluded.

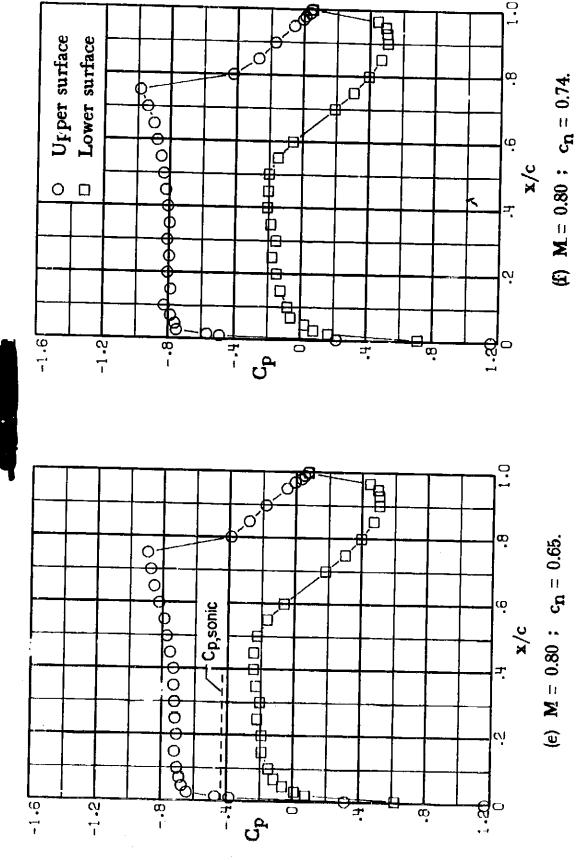


Figure 19. - Continued.

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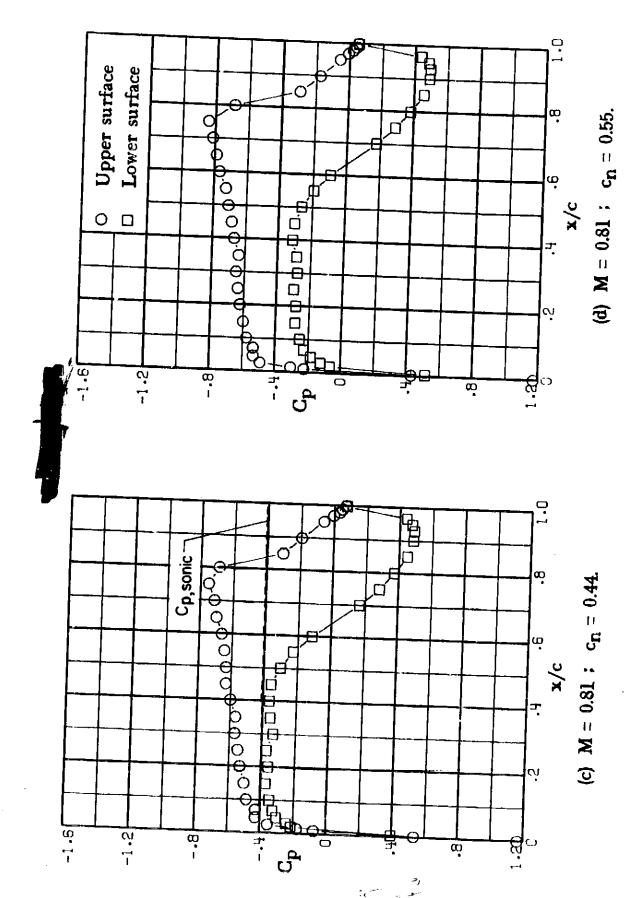


Figure 20. - Continued.

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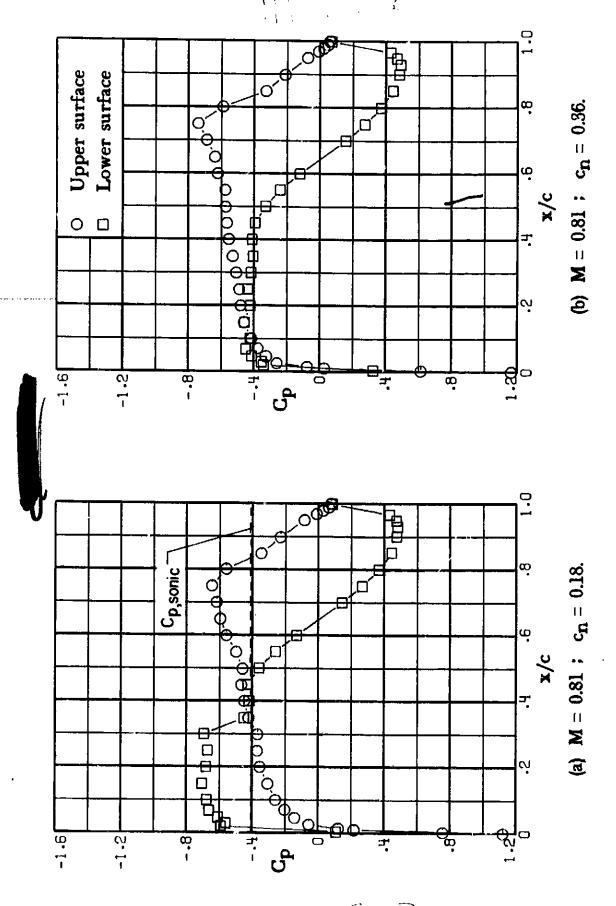
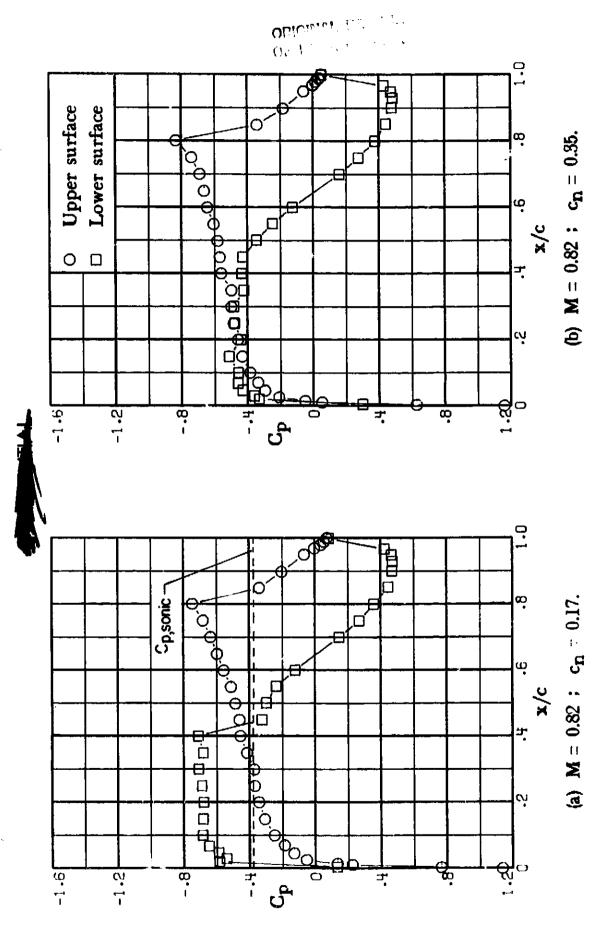


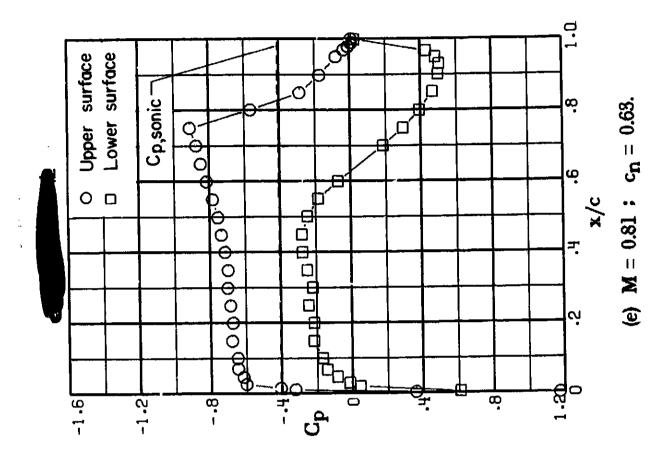
Figure 20. - Chordwise pressure distributions for 10-percent thick supercritical airfoll 33. M=0.81.

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M = 0.82. Figure 21. - Chordwise pressure distributions for 10-percent thick supercritical airfoll 33.

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Figure 20. - Concluded.

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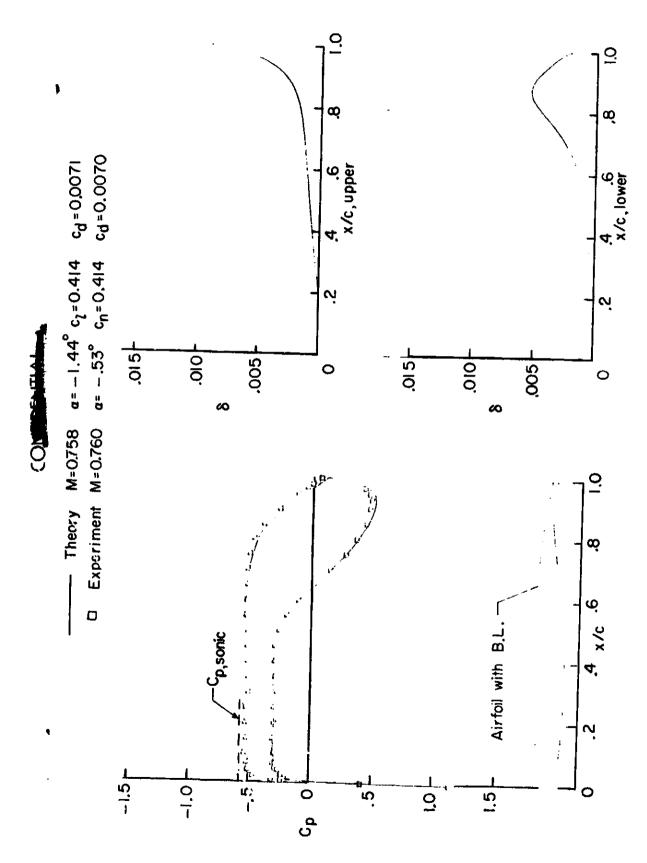


Figure 22. - Comparison of theoretical and experimental characteristics for the 10-percent-thick experimental supercritical airfoil 33. $R=7.7 \times 10^6$.

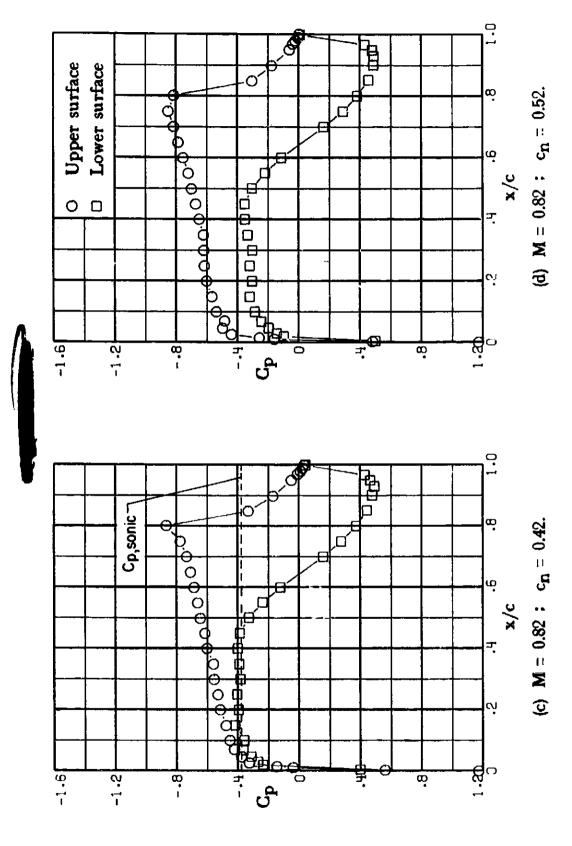


Figure 21. - Concluded.

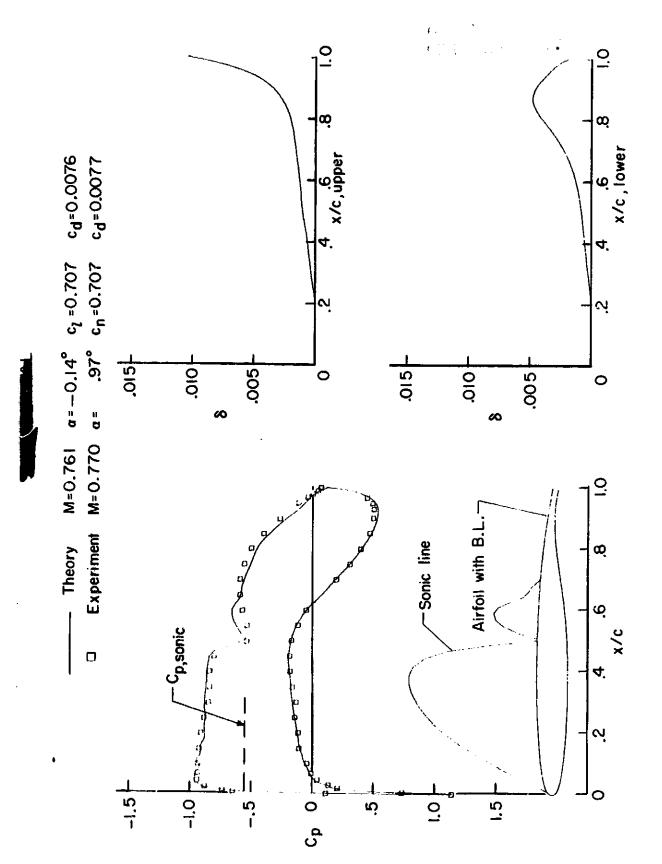


Figure 24. - Comparison of theoretical and experimental characteristics for the 10-percent-thick experimental supercritical airfoil 33. $R = 7.7 \times 10^6$.

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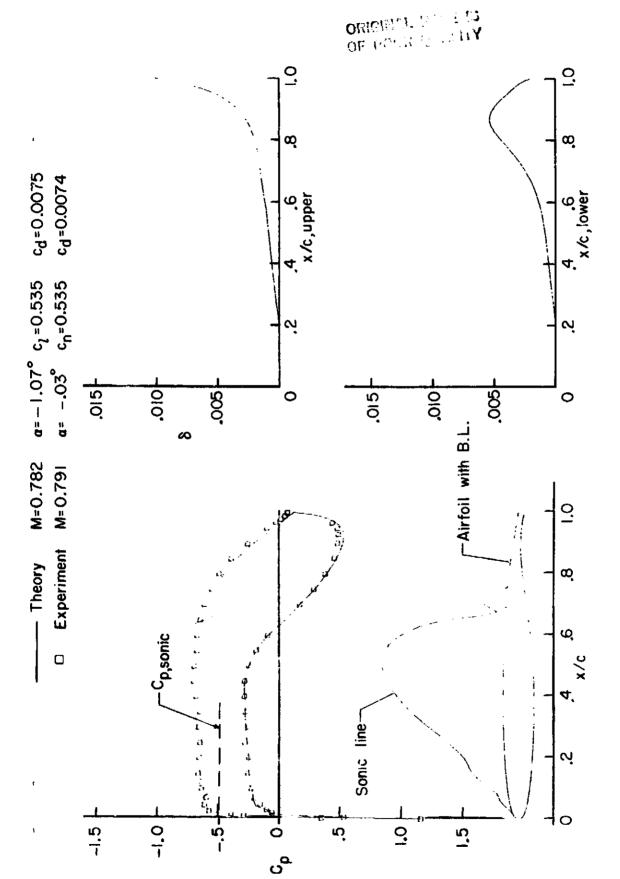


Figure 23. - Comparison of theoretical and experimental characteristics for the 10-percent-thick experimental supercritical airfoil 33. $R = 7.8 \times 10^6$.